# THE NON-CENTRAL $\chi^{2}$ - AND $F$-DISTRIBUTIONS AND THEIR APPLICATIONS $\dagger$ 

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## 1. Introductory

In the Neyman-Pearson theory of testing statistical hypotheses, the efficiency of a statistical test is to be judged by its power of detecting departures from the null hypothesis. Thus besides knowing the random sampling distribution of a given statistic $T$ under this hypothesis, say $H_{0}$, it is also necessary to know the distribution of $T$ under admissible hypotheses alternative to $H_{0}$. Hence the power function of the test is obtained. In the case of the wellknown tests using $\chi^{9}, t$ and $F$, the evaluation of their power functions involves the use of what have been called non-central distributions. For example, if we are applying the $t$-test to examine if a sample has come from a normal population with mean $\mu=0\left(H_{0}\right)$, we know that under $H_{0}, t$ has a $5 \%$ chance of exceeding the $5 \%$ point of its distribution. But in order to compute the power of the test we wish to know the chence that $t$ exceeds this point when $\mu$ has alternative values, not equal to zero. This chance is given by the non-central $t$-integral. This distribution has been studied by Fisher (1931), Neyman (1935), Neyman \& Tokarska (1936) and Johnson \& Welch (1939). In a similar way, the non-central $\chi^{2}$ - and $F$-distributions arise in consideration of the power functions of the $\chi^{2}$ - and variance-ratio tests.

The power function may be used either to determine the extent of the departures from $H_{0}$ in a given direction, which will be detected as significant (at a prescribed level) with a given probability, or it may be used to determine in advance the size of experiment necessary to ensure that a worth-while difference will be established as significant, if it exists. But apart from its value in this connexion, the study of non-central distributions is of considerable interest. The mathematical forms of these distributions of $t, \chi^{2}$ and $F$ have been long known, but their use without extensive tabling has not been easy. The present paper is therefore concerned with two lines of investigation:
(a) The derivation of certain approximations to the probability integrals of (i) non-central $\chi^{2}$, and (ii) the ratio of non-central $\chi^{2}$ to an independent central $\chi^{2}$, which we have termed noncentral $F$. These approximations, depending on tabled functions, permit easy calculation.
(b) Discussion of the ways in which these distributions may be used in connexion with the power functions of statistical tests.

## 2. The non-central $\chi^{2}$-distribution

### 2.1. Geometrical derivation

As is well known, the statistic $\chi^{2}$ is defined as the sum of squares of (say) $n$ independent random deviates, $\xi_{i}$, all drawn from a normal population with mean, 0 , and standard deviation, $\sigma$, viz.

$$
\chi^{2}=\sum_{i=1}^{n} \xi_{i}^{2} / \sigma^{2}
$$

$\dagger$ Part of a thesis approved for the degree of Ph.D. of the Univarsity of London.

If, however, the mean $\xi_{i}$ is $a_{i}$ and we write

$$
x_{i}=\xi_{i}-a_{i},
$$

then we have the non-central $\chi^{2}$ defined by

$$
\chi^{\prime 2}=\sum_{i=1}^{n}\left(x_{i}+a_{i}\right)^{2} / \sigma^{2}
$$

The probability distribution of $\mathcal{X}^{\prime 2}$ has been obtained by Fisher (1928) as a particular case of the distribution of the multiple correlation coefficient. A purely analytical proof was given by Tang (1938). As $\chi^{\prime 2}$ is a generalized form of $\chi^{2}$ it may be of interest to compare ite geometrical representation with the familiar geometry of $\chi^{2}$. We therefore give a direct geometrical derivation of the $\chi^{\prime 1}$-distribution.

Without loss of generality we shall assume in what follows that $\sigma=1$, so that the probability law of $x$ is given by

Then

$$
\begin{equation*}
p(x)=\frac{1}{\sqrt{(2 \pi)}} e^{-\mid x^{4}} \tag{1}
\end{equation*}
$$

$$
\chi^{\prime 2}=\sum_{i=1}^{n} \xi_{i}^{2}
$$

In the $n$-dimensional space of the $\xi$ 's, suppose $O$ is the origin, $P$ the point $\left(\xi_{1}, \ldots, \xi_{n}\right)$, $A$ the point ( $a_{1}, \ldots, a_{n}$ ), $\angle P O A=\theta$ and $M$ the foot of the perpendicular from $P$ on $O A$ as shown in Fig. 1. Then

$$
O P^{2}=\chi^{\prime 2}, \quad O A^{2}=\sum_{i=1}^{n} a_{i}^{2}=\lambda, \quad \text { say } .
$$



Fig. 1
From (1), the probability density at $P$ is proportional to

$$
\begin{equation*}
\exp \left[-\frac{1}{2} \sum_{i=1}^{n}\left(\xi_{i}-a_{i}\right)^{2}\right]=\exp \left[-\frac{1}{2} P A^{2}\right]=\exp \left[-\frac{1}{2}\left(\chi^{\prime 2}+\lambda-2 \chi^{\prime} \sqrt{\lambda \cos \theta)]}\right.\right. \tag{2}
\end{equation*}
$$

If we keep $O P$ and $\theta$ fixed, $P$ describes an $(n-1)$-dimensional sphere of radius $P M=\chi^{\prime} \sin \theta$ with its surface area proportional to $\left(\chi^{\prime} \sin \theta\right)^{n-2}$. If $\chi^{\prime}$ is increased to $\chi^{\prime}+d \chi^{\prime}$ and $\theta$ to $\theta+d \theta$, then a disk of area $\chi^{\prime} d \chi^{\prime} d \theta$ moves round this surface and hence covers a volume proportional to

$$
\left(\chi^{\prime} \sin \theta\right)^{n-8} \chi^{\prime} d \chi^{\prime} d \theta
$$

To obtain the distribution of $\chi^{\prime}$ alone, we integrate out $\theta$. Thus

$$
\left.p\left(\chi^{\prime}\right) d \chi^{\prime}=C \int_{0}^{\pi} e^{-i\left(\chi^{\prime 2}+\lambda-2 \chi^{\prime}\right.} \lambda \cos \theta\right)\left(\chi^{\prime} \sin \theta\right)^{n-2} \chi^{\prime} d \theta d \chi^{\prime},
$$

which is equivalent to

$$
\begin{equation*}
p\left(\chi^{\prime 2}\right) d \chi^{\prime 2}=\frac{C}{2} e^{-\lambda\left(x^{2}+\lambda\right)}\left(\chi^{\prime 8}\right)^{i n-1} d \chi^{\prime 2} \times \int_{0}^{\frac{1 \pi}{n}}\left(e^{-\sqrt{ } \lambda x^{\prime} \cos \theta}+e^{\sqrt{ } \lambda x^{\prime} \cos \theta}\right) \sin ^{n-8} \theta d \theta . \tag{3}
\end{equation*}
$$

Expanding the integrand and integrating term by term, we find

$$
p\left(\chi^{\prime 2}\right)=\frac{1}{2} C e^{-\lambda\left(\chi^{2}+\lambda\right)}\left(\chi^{\prime 2}\right)^{i n-1} \frac{\Gamma\left(\frac{1}{2}\right) \Gamma\left[\frac{1}{2}(n-1)\right]}{\Gamma\left(\frac{1}{2} n\right)}\left\{1+\frac{1}{n}\left(\frac{\chi^{\prime 2} \lambda}{2}\right)+\frac{1}{n(n+2) 2!}\left(\frac{\chi^{\prime 2} \lambda}{2}\right)^{2}+\ldots\right\} .
$$

If zero is substituted for $\lambda$, this reduces to the ordinary $\boldsymbol{\chi}^{2}$-distribution which therefore gives us the value of $C$.

We then have

$$
\begin{equation*}
p\left(\chi^{\prime 2}\right)=\frac{e^{-\dagger x^{2}} e^{-+\lambda}}{2^{i n}} \sum_{j=0}^{\infty} \frac{\left(\chi^{\prime 2}\right)^{)^{n+1}+1} \lambda^{j}}{\Gamma\left(\frac{1}{2} n+j\right)^{2 j} \cdot j!} . \tag{4}
\end{equation*}
$$

### 2.2. Derivation through a transformation of variates

Next we will show that it is possible to effect a variate transformation so as to transform $\chi^{\prime 2}$ into a sum of $(n-1)$ central squares and a single non-central square and then derive its distribution. Make the following orthogonal transformation:

$$
\left.\begin{array}{l}
y_{1}=c_{11} \xi_{1}+c_{18} \xi_{2}+\ldots+c_{1 n} \xi_{n},  \tag{5}\\
\ldots \ldots \ldots \ldots \ldots \ldots \ldots \ldots \ldots \ldots \ldots \ldots \ldots \\
y_{n}=c_{n 1} \xi_{1}+c_{n 2} \xi_{2}+\ldots+c_{n n} \xi_{n} .
\end{array}\right\}
$$

Then

$$
\sum_{1}^{n} \xi_{i}^{q}=\sum_{1}^{n} y_{j}^{2} .
$$

Generally, if

$$
\mathcal{E}\left(y_{j}\right)=c_{j 1} a_{1}+c_{j 2} a_{2}+\ldots+c_{j n} a_{n}=b_{j} \quad(j=1 \text { to } n),
$$

we have

$$
\begin{equation*}
\sum_{1}^{n} a_{i}^{2}=\sum_{1}^{n} b_{j}^{2}, \tag{6}
\end{equation*}
$$

and

$$
\sum_{1}^{n} a_{i} \xi_{i}=\sum_{1}^{n} b_{j} y_{j} .
$$

Now we can make

$$
b_{1}=b_{2}=\ldots=b_{n-1}=0 \quad \text { and } \quad b_{n}=\sqrt{ }\left(\Sigma a_{i}^{2}\right)=\sqrt{ } \lambda .
$$

Thus $\chi^{\prime s}=\sum_{1}^{n} \xi_{i}^{q}$ is distributed as $\sum_{1}^{n-1} y_{j}^{\boldsymbol{q}}+y_{n}^{\mathbf{8}}$, the sum of the squares of $(n-1)$ normal variates with mean zero and the square of a single normal variate with mean $\sqrt{\lambda}$, the s.d.'s being unity.

Writing

$$
\chi^{\prime 2}=w, \quad \sum_{1}^{n-1} y_{j}^{2}=u \quad \text { and } \quad y_{n}^{2}=v,
$$

we see that $u$ has a $\chi^{2}$-distribution with ( $n-1$ ) degrees of freedom, that is,
and that $v$ follows the law

$$
\begin{aligned}
p(v) & =\frac{v^{-t}}{2 \sqrt{ }(2 \pi)}\left\{e^{-H(v v-\sqrt{ })^{2}}+e^{-H(-\sqrt{ } v-\sqrt{ } \lambda)^{2}}\right\} \\
& =\frac{1}{2^{4} \Gamma\left(\frac{1}{2}\right)} e^{-H(v+\lambda)} v^{-t}\left(1+\frac{v \lambda}{2!}+\frac{(v \lambda)^{2}}{4!}+\ldots\right) .
\end{aligned}
$$

Hence, replacing $v$ by $(w-u)$ in the joint probability law $p(u, v)$, we have

$$
p(u, w)=\frac{e^{-1 w} e^{-\dagger \lambda} w^{1(n-1)}}{2^{1 n} \Gamma\left(\frac{1}{2}\right) \Gamma\left[\frac{1}{2}(n-1)\right]}\left(\frac{u}{w}\right)^{1(n-3)}\left\{\left(1-\frac{u}{w}\right)^{-i}+\frac{w \lambda}{2!}\left(1-\frac{u}{w}\right)^{\dagger}+\ldots\right\} .
$$

Whence integrating with respect to $u$ from 0 to $w$, we obtain
that is, $\quad p\left(\chi^{\prime 2}\right)=\frac{e^{-\downarrow \chi^{\prime 2}} e^{-1 \lambda}\left(\chi^{\prime 8}\right)^{\dagger n-1}}{2^{\operatorname{tn}} \Gamma\left(\frac{1}{2} n\right)}\left\{1+\frac{1}{n}\left(\frac{\chi^{\prime 2} \lambda}{2}\right)+\frac{1}{n(n+2) .2!}\left(\frac{\chi^{\prime 2} \lambda}{2}\right)^{8}+\ldots\right\}$,
Which is seen to be the same as (4).
In this distribution of $\chi^{\prime 2}, n$ may be called the number of degrees of freedom and $\lambda$, which is equal to the sum of the squares $\sum_{1}^{n} a_{i}^{i}$, the non-central parameter.

### 2.3. Conditional distribution of $\chi^{\prime 8}$ under linear constraints

Suppose the $\xi$ 's are subject to $k(<n)$ linear constraints. These can be transformed into an orthogonal set represented, say, by the equations
where

$$
\begin{equation*}
\sum_{i=1}^{n} c_{j i} \xi_{i}=\rho_{j} \quad(j=1, \ldots, k) \tag{8}
\end{equation*}
$$

$$
\sum_{i=1}^{n} c_{j i}^{2}=1, \quad \sum_{i=1}^{n} c_{j i} c_{i i}=0 \quad(j \neq l)
$$

We make an orthogonal transformation of variates defined by the equations (5), so that $\Sigma \xi_{i}^{q}$ transforms to $\Sigma y_{j}^{2}$ and the $k$ constraints of (8) become simply $y_{1}=\rho_{1}, \ldots, y_{k}=\rho_{k}$. To find the distribution of $\Sigma y_{j}^{\ell}$ subjeot to these conditions, we first see that, in virtue of the relations in (6), the joint probability law of the $\xi$ 's
transforms into

$$
\begin{aligned}
& p\left(\xi_{1}, \ldots, \xi_{n}\right)=C \exp \left\{-\frac{1}{2} \sum_{1}^{n}\left(\xi_{i}-a_{i}\right)^{2}\right\} \\
& p\left(y_{1}, \ldots, y_{n}\right)=C \exp \left\{-\frac{1}{2} \sum_{1}^{n}\left(y_{j}-b_{j}\right)^{2}\right\}
\end{aligned}
$$

When $y_{1}, \ldots, y_{k}$ take respectively the constant values $\rho_{1}, \ldots, \rho_{k}$, we have the conditional probability law

$$
\begin{equation*}
p\left(y_{k+1}, \ldots, y_{n} \mid \rho_{1}, \ldots, \rho_{k}\right)=C_{1} \exp \left\{-\frac{1}{2} \sum_{k+1}^{n}\left(y_{j}-b_{j}\right)^{2}\right\} . \tag{9}
\end{equation*}
$$

It can be shown from (9), as in $\S 2 \cdot 1$, that the sum of the non-central squares $\left(y_{k+1}^{*}+\ldots+y_{n}^{2}\right)$ is distributed as a $\chi^{\prime 2}$ with $(n-k)$ degrees of freedom and parameter

$$
\lambda=b_{k+1}^{2}+\ldots+b_{n}^{2}
$$

From (6) we see that
and

$$
\begin{align*}
y_{k+1}^{2}+\ldots+y_{n}^{8} & =\Sigma \xi_{i}^{2}-\left(\rho_{i}^{2}+\ldots+\rho_{k}^{2}\right) \\
b_{k+1}^{\ell}+\ldots+b_{n}^{2} & =\Sigma a_{i}^{i}-\left(b_{1}^{8}+\ldots+b_{k}^{2}\right) \\
& =\sum_{i=1}^{n} a_{i}^{2}-\sum_{j=1}^{k}\left(\sum_{i=1}^{n} a_{i} c_{j i}\right)^{2} . \tag{10}
\end{align*}
$$

Hence $\left(\sum_{1}^{n} \xi_{i}^{2}-\sum_{1}^{k} \rho_{3}^{9}\right)$ is distributed as a $\chi^{\prime 2}$ with $(n-k)$ degrees of freedom and parameter $\lambda$, given by the expression in (10).

In particular, if there is only a single constraint on the $\xi$ 's, given by

$$
\begin{equation*}
\sum_{1}^{n} c_{i} \xi_{i}=\rho, \quad \sum_{1}^{n} c_{i}^{i}=1, \tag{11}
\end{equation*}
$$

then $\left(\sum_{1}^{n} \xi_{i}^{q}-\rho^{2}\right)$ follows a $\chi^{\prime 2}$-distribution with $(n-1)$ degrees of freedom and

$$
\begin{equation*}
\lambda=\sum_{1}^{n} a_{i}^{i}-\left(\sum_{1}^{n} a_{i} c_{i}\right)^{2} . \tag{12}
\end{equation*}
$$

## 3. Approximations to the $\chi^{\prime 2}$-distribution

### 3.1. The $\chi^{2}$-approximation

Fisher (1928) has shown that the distribution function of $\chi^{\prime 8}$ given by (4) can be expressed in terma of a Bessel function with imaginary argument. When $n$, the number of degrees of freedom, is odd, this can be reduced to elementary functions. When $n$ is even, we see that the probability integral

$$
\int_{\chi^{n}}^{\infty} p\left(\chi^{\prime 2}\right) d \chi^{\prime 2}
$$

can be expressed as a double Poisson sum. However, in both cases, the labour of calculating the probability integral is considerable.

In his paper, Fisher has given a table of the upper $5 \%$ significance points of the $\chi^{\prime 1}$ distribation for $n=1$ to 7 and $\sqrt{\lambda}=0(0 \cdot 2) 5 \cdot 0$. Garwood $\dagger$ has an unpublished table of the lower $5 \%$ points for the same range of values of $n$ and $\lambda$. No tables of the probability integral are available. It may therefore be useful to have an easy method of determining the probability integral and percentage points sufficiently accurately for any given values. For this purpose we shall consider several approximations to the distribution of $\chi^{\prime 2}$.

The characteristio function of this distribution is easily seen to be

$$
\phi(t)=\exp \left\{\frac{\lambda i t}{1-2 i t}\right\} /(1-2 i t)^{+n} .
$$

Hence we have the following oumulants:

$$
\left.\begin{array}{ll}
\kappa_{1}=n+\lambda, & \kappa_{2}=2(n+2 \lambda),  \tag{13}\\
\kappa_{3}=8(n+3 \lambda), & \kappa_{4}=48(n+4 \lambda),
\end{array}\right\}
$$

the general $r$ th cumulant being

$$
\kappa_{r}=2^{r-1}(r-1)!(n+r \lambda) .
$$

In the $\beta_{1}, \beta_{2}$ diagram, it was found that the point computed from the above $\kappa$ 's moved close to and above the Type III line, and this auggested that we might fit a Type III distribution from the first two moments. This is given by

$$
\begin{equation*}
f(y)=\frac{e^{-t y} y^{\downarrow v-1}}{2^{\downarrow \nu} \Gamma\left(\frac{1}{2} \nu\right)}, \tag{14}
\end{equation*}
$$

where

$$
y=\chi^{\prime 2} / \rho,
$$

$$
\begin{equation*}
\rho=\frac{n+2 \lambda}{n+\lambda}=1+\frac{\lambda}{n+\lambda}, \quad \nu=\frac{(n+\lambda)^{2}}{(n+2 \lambda)}=n+\frac{\lambda^{2}}{n+2 \lambda} . \tag{15}
\end{equation*}
$$

This means that we are representing the distribution of $\left(\chi^{\prime 2} / \rho\right)$ by that of $\chi^{2}$ with $\nu$ degrees of freedom, $\nu$ being in general a fraction.
$\dagger$ I am grateful to Dr F. Garwood for kindly making his table available to me for reference.

In what follows we shall write $x$ for $\chi^{\prime 2}, p(x)$ for the true distribution of $\chi^{\prime 2}$ with $n$ degrees of freedom and parameter $\lambda$ and $f(x)$ for the approximation to $p(x)$ obtained by assuming that $x / \rho=y$ is distributed as $\chi^{2}$ with $\nu$ degrees of freedom.

Then the probability integral

$$
\int_{0}^{x} p(x) d x=\int_{0}^{y} p(y) d y
$$

is approximately given by

$$
\int_{0}^{y} f(y) d y
$$

This integral can be expressed in the notation of the tables of the Incomplete $\Gamma$-function (K. Pearson, 1922) as $I(u, p)$, where

$$
\begin{equation*}
u=\frac{y}{\sqrt{(2 \nu)}}=\frac{x}{\sqrt{[2(n+2 \lambda)]}}, \quad p=\frac{\nu}{2}-1=\frac{(n+\lambda)^{2}}{2(n+2 \lambda)}-1 \tag{16}
\end{equation*}
$$

and could be evaluated by interpolation in these tables. For interpolation $u$-wise the second differences with Everett interpolation coefficients may be used, while linear interpolation $p$-wise seems adequate.

The approximations to the probability integral so obtained for certain values of $n, \lambda$ and $x$ are shown in Table 1 for comparison with the exact values. In some of these cases $x$ is the upper $5 \%$ point (Fisher) or the lower $5 \%$ point (Garwood), so that the exact values are 0.95 or 0.05 . The others are directly computed. For many purposes, especially in connexion with power functions, the degree of accuracy given by this method may be considered quite adequate.

Table 1. Showing exact and approximate values of the $\chi^{\prime 2}$-integrals, $\int_{0}^{x} p(x) d x$

| $n$ | $\lambda$ | $x$ | Approx. | Exact |
| :---: | :---: | :---: | :---: | :---: |
| 4 | $\begin{array}{r} 4 \\ 4 \\ 4 \\ 4 \\ 10 \end{array}$ | $\begin{array}{r} 1.765 \\ 10.000 \\ 17.309 \\ 24.000 \\ 10.000 \end{array}$ | $\begin{aligned} & 0.0399 \\ & 0.711 \\ & 0.9492 \\ & 0.9913 \\ & 0.3178 \end{aligned}$ | $\begin{aligned} & 0.0500 \\ & 0.7118 \\ & 0.9500 \\ & 0.9925 \\ & 0.3148 \end{aligned}$ |
| 7 | $\begin{array}{r} 1 \\ 1 \\ 16 \\ 16 \\ 16 \end{array}$ | 4.000 16.004 10.257 24.000 38.970 | $\begin{aligned} & 0.1621 \\ & 0.9499 \\ & 0.0430 \\ & 0.5947 \\ & 0.0482 \end{aligned}$ | $\begin{aligned} & 0.1628 \\ & 0.9500 \\ & 0.0500 \\ & 0.5888 \\ & 0.9500 \end{aligned}$ |
| 12 | $\begin{array}{r} 6 \\ 18 \end{array}$ | $\begin{aligned} & 24 \cdot 000 \\ & 24 \cdot 000 \end{aligned}$ | $\begin{aligned} & 0.8187 \\ & 0.2936 \end{aligned}$ | $\begin{aligned} & 0.8174 \\ & 0.2901 \end{aligned}$ |
| 16 | $\begin{array}{r} 8 \\ 8 \\ \mathbf{8 2} \\ 32 \end{array}$ | 30.000 <br> 40.000 <br> $30 \cdot 000$ <br> 60.000 | $\begin{aligned} & 0.7896 \\ & 0.9682 \\ & 0.0590 \\ & 0.8329 \end{aligned}$ | $\begin{aligned} & 0.7880 \\ & 0.9632 \\ & 0.0609 \\ & 0.8318 \end{aligned}$ |
| 24 | $\begin{aligned} & 24 \\ & 24 \\ & 24 \end{aligned}$ | $\begin{aligned} & 36 \cdot 000 \\ & 48 \cdot 000 \\ & 72.000 \end{aligned}$ | $\begin{aligned} & 0.1556 \\ & 0.5333 \\ & 0.9658 \end{aligned}$ | $\begin{aligned} & 0.1567 \\ & 0.5298 \\ & 0.8667 \end{aligned}$ |

To find the percentage points of the $\chi^{\prime 2}$ distribution, we first interpolate in the appropriate percentage point tables of the $\chi^{2}$ (e.g. Thompson, 1941) for $\nu$ degrees of freedom and then multiply the interpolate by $\rho$. Four-point Lagrangian interpolation formulae may be used. The approximate upper and lower $5 \%$ points obtained by this method for certain values of $n$ and $\lambda$ are given in Table 2, along with the exaot values. Clearly the accuracy is not as good fer the lower points as for the upper ones. Although the comparisons have had to be confined only to small values of $n$, since Fisher and Garwood have only given exact percentage points up to $n=7$, from the closeness of the probability integral approximation (Table 1) we could still expect that the approximation to the percentage points would be fairly close for higher $n$.

These approximations based on the $\chi^{2}$ fit will be referred in subsequent sections as the first approximation.

Table 2. Showing exact and approximate values of the percentage points of the $\chi^{\prime 2}$-distribution

| $n$ | $\lambda$ | Upper $5 \%$ point |  | Lower 5\% point |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | Approx. | Exact | Approx. | Exaot |
| 2 | 1 4 16 25 | $\begin{array}{r} 8 \cdot 63 \\ 14 \cdot 72 \\ 33 \cdot 35 \\ 45 \cdot 66 \end{array}$ | $\begin{array}{r} 8 \cdot 64 \\ 14 \cdot 64 \\ 33 \cdot 06 \\ 45 \cdot 31 \end{array}$ | $\begin{array}{r} 0.20 \\ 0.94 \\ 6.89 \\ 12.68 \end{array}$ | $\begin{array}{r} 0.17 \\ 0.65 \\ 6.32 \\ 12.08 \end{array}$ |
| 4 | 1 4 16 25 | $\begin{aligned} & 11 \cdot 72 \\ & 17.38 \\ & 35 \cdot 69 \\ & 47.94 \end{aligned}$ | $\begin{aligned} & 11 \cdot 71 \\ & 17 \cdot 31 \\ & 35 \cdot 43 \\ & 47 \cdot 61 \end{aligned}$ | $\begin{array}{r} 0.93 \\ 1.95 \\ 8.36 \\ 14.26 \end{array}$ | $\begin{array}{r} 0.91 \\ 1.77 \\ 7.88 \\ 13.73 \end{array}$ |
| 7 | 1 4 16 25 | $\begin{aligned} & 16.01 \\ & 21.28 \\ & 39.16 \\ & 51.34 \end{aligned}$ | $\begin{aligned} & 16.00 \\ & 21.23 \\ & 38.07 \\ & 51.06 \end{aligned}$ | $\begin{array}{r} 2.51 \\ 3.78 \\ 10.64 \\ 16.68 \end{array}$ | $\begin{array}{r} 2.49 \\ 3.66 \\ 10.26 \\ 16.23 \end{array}$ |

### 3.2. The normal approximation

It is known that, for $n>30$, Fisher's approximation, that $\sqrt{ }\left(2 \chi^{2}\right)$ is distributed as a normal variate $N(\sqrt{ }(2 n-1), 1), \dagger$ will give fairly close values to the probability integral and percentage points of the $\chi^{2}$-distribution. It can be shown that a similar normal approximation is available for the $\chi^{\prime 2}$-distribution for large values of $n$ or $\lambda$.

First we shall show that $\chi^{\prime}$ approaches normality with greater rapidity than $\chi^{\prime 2}$.
If $x$ is written for $\chi^{\prime 2}$, and $x_{0}$ is mean $x$, we have by Taylor's theorem

$$
\begin{aligned}
& x^{\sharp}=x_{0}^{\ddagger}+\frac{1}{8}\left(x-x_{0}\right) x_{0}^{-\frac{1}{4}}-\frac{1}{8}\left(x-x_{0}\right)^{2} x_{0}^{-\frac{4}{4}}+\frac{1}{18}\left(x-x_{0}\right)^{3} x_{0}^{-4}+\ldots, \\
& x^{4}=x_{0}^{\frac{1}{8}}+\frac{3}{8}\left(x-x_{0}\right) x_{0}^{\frac{j}{3}}+\frac{3}{8}\left(x-x_{0}\right)^{2} x_{0}^{-\frac{1}{2}}-\frac{1}{18}\left(x-x_{0}\right)^{3} x_{0}^{-4}+\ldots .
\end{aligned}
$$

[^0]By taking expectations on both sides and substituting from (13) the moments of $x=\chi^{\prime 8}$, we get $\mu_{1}^{\prime}$ and $\mu_{3}^{\prime}$ of $\chi^{\prime}$. Also

$$
\mu_{2}^{\prime}\left(\chi^{\prime}\right)=\mu_{1}^{\prime}\left(\chi^{\prime 8}\right), \quad \mu_{4}^{\prime}\left(\chi^{\prime}\right)=\mu_{2}^{\prime}\left(\chi^{\prime 2}\right)
$$

Hence we derive the following moments:

$$
\begin{aligned}
& \mu_{1}^{\prime}=(n+\lambda)^{t}-\frac{n+2 \lambda}{4(n+\lambda)^{t}}+\frac{1}{2} \frac{n+3 \lambda}{(n+\lambda)^{t}}-\frac{15}{32} \frac{(n+2 \lambda)^{2}}{(n+\lambda)^{\frac{1}{4}}}+\ldots, \\
& \mu_{2}^{\prime}=(n+\lambda), \\
& \mu_{3}^{\prime}=(n+\lambda)^{t}+\frac{3}{4} \frac{n+2 \lambda}{(n+\lambda)^{4}}-\frac{1}{2} \frac{n+3 \lambda}{(n+\lambda)^{t}}+\frac{9}{32} \frac{(n+2 \lambda)^{2}}{(n+\lambda)^{\frac{1}{4}}}+\ldots, \\
& \mu_{4}^{\prime}=2(n+2 \lambda)+(n+\lambda)^{2}
\end{aligned}
$$

from which we obtain

Hence

$$
\begin{array}{rlr}
\mu_{1}^{\prime}=(n+\lambda)^{1}-\frac{n+2 \lambda}{4(n+\lambda)^{4}}+\ldots, \quad \mu_{2}=\frac{n+2 \lambda}{2(n+\lambda)}+\ldots, \\
\mu_{3} & =\frac{n+3 \lambda}{(n+\lambda)^{4}}-\frac{3}{4} \frac{(n+2 \lambda)^{2}}{(n+\lambda)^{4}}+\ldots, \quad \mu_{4}=\frac{3}{4}+O\left[(n+\lambda)^{-2}\right] . \\
\gamma_{1}=\frac{\mu_{3}}{\mu_{2}}=\frac{n^{2}+4 n \lambda}{\sqrt{2(n+\lambda)(n+2 \lambda)^{1}}+\ldots,} \quad \gamma_{2}=\frac{\mu_{4}^{4}-3=O\left[(n+\lambda)^{-2}\right] .}{\mu_{2}^{9}} .
\end{array}
$$

Comparing these with the corresponding coefficients of the $\chi^{\prime 2}$-distribution, viz.

$$
\gamma_{1}=\frac{\sqrt{8(n+3 \lambda)}}{(n+2 \lambda)^{4}}+\ldots, \quad \gamma_{2}=\frac{12(n+4 \lambda)}{(n+2 \lambda)^{2}}+\ldots
$$

we see that $\chi^{\prime}$ approaches normality faster than $\chi^{\prime 2}$.
From the above it follows that $\sqrt{ }\left(2 \chi^{2}\right)$ has mean $\sqrt{ }\{2(n+\lambda)-(n+2 \lambda) /(n+\lambda)\}$ to order $(n+\lambda)^{-1}$ and variance $(n+2 \lambda) /(n+\lambda)$ to order $(n+\lambda)^{-1}$. We can therefore regard

$$
\sqrt{\left\{\frac{2 \chi^{\prime 2}(n+\lambda)}{n+2 \lambda}\right\}}
$$

as distributed normally with mean

$$
\sqrt{\left\{\frac{2(n+\lambda)^{2}}{n+2 \lambda}-1\right\}}
$$

and variance unity.
This result may also be derived by taking the $\chi^{8}$-approximation to the $\chi^{\prime 2}$-distribution and then using the known result that for large $\nu, \sqrt{ }\left(2 \chi^{3}\right)$ is distributed as $N[\sqrt{ }(2 \nu-1), 1]$. For, substituting $\chi^{\prime 8} / \rho$ for $\chi^{2}$ and the expressions in (15) for $\rho$ and $\nu$, we reach the same normal approximation.

Since $\nu>n$ from (15), it can be seen that the normal approximation to $\chi^{\prime}$ with $n$ degrees of freedom will be better than the normal approximation to $\chi$ with the same degrees of freedom. Thus, for example, if $n=25$, we have

$$
\begin{array}{ccccc}
\lambda=0 & 10 & 20 & 30 & 40 \\
\nu=25 & 27 \cdot 22 & 31 \cdot 15 & 35 \cdot 59 & 40 \cdot 24 .
\end{array}
$$

Hence for sufficiently large values of $n$ and $\lambda$, the probability integral and percentage points may be obtained from the normal tables. Table 3 gives a comparison of some values of the probability integral, thus calculated, with the exact values.

Table 3. Values of the $\chi^{\prime 2}$-integral on the normal approximation

| $n$ | $\lambda$ | $\nu$ | $x$ | From <br> $\chi^{2}$ | From <br> normal | Exact |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 16 | 32 | 28.8 | 30 | 0.0590 | 0.0638 | 0.0809 |
| 16 | 32 | 28.8 | 60 | 0.8329 | 0.8320 | 0.8316 |
| 24 | 24 | 32.0 | 36 | 0.1556 | 0.1515 | 0.1567 |
| $\mathbf{2 4}$ | 24 | 32.0 | 72 | 0.9656 | 0.9680 | 0.9867 |

3.3. Closer approximations to the $\chi^{\prime 2}$-distribution

The probability function of $\chi^{\prime 9}$ can be represented in the form of a series with the fitted probability function of ( $\rho \chi^{2}$ ) as the leading term and, from these mathematical expansions, closer approximations to the probability integral and percentage points may be obtained. Two methods will be briefly considered.

First method
The cumulants of the distribution $f(x)$, as defined on p. 207 above, are seen to be

$$
\left.\begin{array}{cc}
\kappa_{1}^{*}=n+\lambda, & \kappa_{2}^{*}=2(n+2 \lambda),  \tag{17}\\
\kappa_{3}^{*}=\frac{8(n+2 \lambda)^{2}}{n+\lambda}, & \kappa_{4}^{*}=\frac{48(n+2 \lambda)^{3}}{(n+\lambda)^{2}}
\end{array}\right\}
$$

the rth cumulant being

$$
\begin{equation*}
\kappa_{3}-\kappa_{3}^{*}=c_{3}, \quad \kappa_{4}-\kappa_{4}^{*}=c_{4}, \ldots \tag{18}
\end{equation*}
$$

Then the corresponding differences of cumulants of $p(y)$ and $f(y)$ as defined on $p, 207$, will be

$$
c_{3} / \rho^{3}, \quad c_{4} / \rho^{4}, \ldots
$$

By the application of the Edgeworth operator to $f(y)$ we have

$$
\begin{aligned}
p(y) & =\exp \left\{-\frac{c_{3} d^{3}}{6 \rho^{3} d y^{3}}+\frac{c_{4} d^{4}}{24 \rho^{4} d y^{4}}+\ldots\right\} f(y) \\
& =\left[1+\left\{-\frac{c_{3}}{6 \rho^{3}} D^{3}+\frac{c_{4}}{24 \rho^{4}} D^{4}-\ldots\right\}+\frac{1}{2!}\left(\left(\frac{c_{3}}{6 \rho^{3}}\right)^{2} D^{6}+\left(\frac{c_{4}}{24 \rho^{4}}\right)^{2} D^{8}-\ldots\right\}+\ldots\right] f(y) .
\end{aligned}
$$

Hence the probability integral $\int_{0}^{y} p(y) d y$ is given by

$$
\begin{align*}
& \int_{0}^{\nu} f(y) d y+\left[\left\{-\frac{c_{3}}{8 \rho^{3}} f^{\prime \prime}(y)+\frac{c_{4}}{24 \rho^{4}} f^{\prime \prime}(y)+\ldots\right\}\right. \\
&  \tag{19}\\
& \left.\quad+\frac{1}{2!}\left\{\left(\frac{c_{3}}{6 \rho^{3}}\right)^{2} f^{(5)}(y)+\left(\frac{c_{4}}{24 \rho_{4}}\right)^{2} f^{(7)}(y)+\ldots\right\}+\ldots\right] .
\end{align*}
$$

Since the higher derivatives of $f(y)$ become smaller in value for a given $y$, we retain only the first term in the square brackets of (19) and get a second approximation to the probability integral in the form

$$
\int_{0}^{y} f(y) d y-\frac{c_{3} d^{3}}{6 \rho^{3} d y^{3}} \int_{0}^{y} f(y) d y
$$

which can be written as

$$
\begin{equation*}
I(u, p)-\frac{c_{3}}{6 \rho^{3}[\sqrt{ }(2 \nu)]^{3} d u^{3}} . \tag{20}
\end{equation*}
$$

When using the expression (20) for the evaluation of the integral, the computation of the first term $I(u, p)$ will, in general, require interpolation in the tables of the Incomplete $\Gamma$ function. We shall now show that by a suitable modification of the Everett interpolation formula, the second term in (20) can be accounted for and the whole expression computed in one calculation.

If $u_{1}, u_{2}$ are the tabulated values between which $u$ lies and $\Delta_{1}^{\prime \prime}, \Delta_{2}^{\prime \prime}$ the tabulated second differences, we have as an approximation

$$
\frac{d^{3} I}{d u^{3}} \sim\left(\Delta_{2}^{\prime \prime}-\Delta_{1}^{\prime \prime}\right) 10^{3},
$$

the interval for $u$ being 0.1 in the tables. Suppose $q$ is the fraction $\left(u-u_{1}\right) /\left(u_{2}-u_{1}\right), E_{1}^{*}, E_{2}^{*}$ the second-order Everett interpolation coefficients corresponding to $q$ and $k=\frac{c_{3} \frac{10^{3}}{6 \rho^{3}[\sqrt{(2 \nu)}]^{3}}}{}$. Then (20) becomes

$$
\begin{equation*}
I\left(u_{1}, p\right)(1-q)+I\left(u_{2}, p\right) q+\Delta_{1}^{\prime \prime}\left(E_{1}^{\prime \prime}+k\right)+\Delta_{2}^{*}\left(E_{2}^{\pi}-k\right) \tag{21}
\end{equation*}
$$

If $p$ is not a tabled value but lies between $p_{1}$ and $p_{2}$, then we evaluate the above expression for $p_{1}$ and for $p_{2}$ and then interpolate linearly for $p$.

## Second method

It is well known that by using the Edgeworth form of the Gram-Charlier Type A series, a frequency function can be normalized if it approaches normality asymptotically and if its cumulants are in increasing order of some quantity, $n^{-1}$.

Goldberg \& Levine (1946) have shown that by the method of normalization the percentage points of the $\chi^{2}$-distribution could be obtained to a fairly good degree of accuracy. A similar method might be applied usefully to the $\chi^{\prime 2}$-distribution. However, a modified form of expansion with the fitted $\chi^{2}$-function as the first term will be found more suitable.

Let us standardize the variate $x$ (written for $\chi^{\prime 8}$ ) by introducing

$$
\xi=\frac{x-(n+\lambda)}{\sqrt{(2 n+4 \lambda)^{\circ}}}
$$

Then, using the same notation as before, the cumulants of the distribution $p(\xi)$ sie

$$
0, \quad 1, \quad \kappa_{3} / \kappa \frac{1}{2}, \quad \kappa_{4} / \kappa_{2}^{\mathbf{3}}, \ldots .
$$

Since $f(x)$ has the same mean and standard deviation as $p(x)$, we get for the cumulants of $f(\xi)$

$$
0, \quad 1, \quad \kappa_{3}^{*} / \kappa \frac{1}{1}, \quad \kappa_{4}^{*} / \kappa_{8}^{2}, \ldots
$$

These cumulants, from the third onwards, are of orders $-\frac{1}{2},-1,-\frac{3}{2}, \ldots$ in both $n$ and $\lambda$. Now let

$$
\alpha=\alpha(\xi)=e^{-k \xi^{2}} / \sqrt{ }(2 \pi),
$$

and let $\xi_{3}, \xi_{4}, \ldots$ be the Hermite polynomials of orders 3, 4, $\ldots$. Then we have, arranging the terms in order of magnitude of $n$ (Kendall, I, 1945, § 6.32),

$$
\begin{equation*}
p(\xi)=\alpha(\xi)+\frac{1 \kappa_{3}}{6 \kappa_{2}^{1}} \alpha \xi_{3}+\alpha\left(\frac{1 \kappa_{4}}{24 \kappa_{2}^{3} \xi_{4}}+\frac{1}{72} \frac{\kappa_{8}^{8}}{\kappa_{2}^{3}} \xi_{\varnothing}\right)+\ldots . \tag{22}
\end{equation*}
$$

There is a similar expansion for $f(\xi)$ with $\kappa_{r}^{*}$ in place of $\kappa_{r}(r>2)$.

Now we subtract formally this second series from the first, term by term, and transfer $f(\xi)$ to the right-hand side. We then obtain

$$
\begin{equation*}
p(\xi)=f(\xi)+\alpha(\xi)\left[\frac{1}{6} \frac{c_{3}}{\kappa \frac{1}{2}} \xi_{3}+\left(\frac{1}{24} \frac{c_{4}}{\kappa_{2}^{3}} \xi_{4}+\frac{1}{72} \frac{c_{33}}{\kappa_{2}^{3}} \xi_{6}\right)+\ldots\right], \tag{23}
\end{equation*}
$$

where $c_{3} c_{4}$ have the same meanings as in (18) and $c_{r h l}$ is written for ( $\kappa_{r} \kappa_{h} \kappa_{l}-\kappa_{r}^{*} \kappa_{h}^{*} \kappa_{l}^{*}$ ).
We know that the infinite series in (23) is not uniformly convergent. We can still integrate it formally term by term and make use of the first few terms to get a better approximation than that given by the integral of $f(\xi)$ alone. Thus retaining terms up to $O\left(n^{-\mathbf{t}}\right)$, we derive an approximation to the probability integral

$$
\int_{0}^{x} p(x) d x=\int_{0}^{\xi} p(\xi) d \xi
$$

in the form

$$
\begin{equation*}
\int_{0}^{\xi} f(\xi) d \xi+\alpha(\xi)\left[-\frac{1}{6} \frac{c_{3}}{\kappa!} \xi_{2}-\left(\frac{1}{24} \frac{c_{4}}{\kappa_{2}^{3}} \xi_{3}+\frac{1}{72} \frac{c_{33}}{\kappa_{2}^{3}} \xi_{5}\right)-\left(-\frac{1}{120} \frac{c_{5}}{\kappa_{2}^{\prime}} \xi_{4}+\frac{1}{144} \frac{c_{34}}{\kappa_{2}^{\frac{1}{2}}} \xi_{6}+\frac{1}{1296} \frac{c_{333}}{\kappa_{8}^{!}} \xi_{8}\right)\right] . \tag{24}
\end{equation*}
$$

The first term in (24) is our first approximation of §3.1 and the rest give a correction to it which is seen to result in a considerable improvement (see Table 4). For evaluating this expression, the values of the Hermite polynomials may be taken from Jorgensen's tables (1916) if $\xi$ is an argument tabled there; otherwise they have to be directly calculated. $\alpha(\xi)$ may be found (without need for interpolation) from Tables of the Probability Functions, Vol. 2 (Federal Works Agency, New York, 1942).
The coefficients in (24) involve only differences of the cumulants and so are smaller than the corresponding coefficients in (22). Thus a closer approximation is likely to result from (24) than from the same order of terms in (22).
For the percentage points, we employ the inversion of the Gram-Charlier series obtained by Cornish \& Fisher (1937). If $x, x^{\prime}$ and $\xi$ are respectively the percentage points of the distributions $p(x), f(x)$ and $\alpha(\xi)$, then for a given probability level, we have

$$
\frac{x-(n+\lambda)}{\sqrt{(2 n+4 \lambda)}}=\xi+\frac{1}{6} \frac{\kappa_{3}}{\kappa_{2}^{\frac{1}{2}}}\left(\xi^{3}-1\right)+\left\{\frac{1}{24} \frac{\kappa_{4}}{\kappa_{2}^{3}}\left(\xi^{3}-2 \xi\right)-\frac{1}{36} \kappa_{8}^{2}\left(2 \xi^{3}-5 \xi\right)\right\}+\ldots
$$

$\frac{x^{\prime}-(n+\lambda)}{\sqrt{(2 n+4 \lambda)}}$ has a similar expansion with $\kappa_{r}^{*}$ in place of $\kappa_{r}(r>2)$. By differencing as before we obtain an expression for $x$ in terms of $x^{\prime}$ and $\xi$. Retaining terms up to $O\left(n^{-\mathbf{I}}\right)$, we find

$$
\begin{align*}
x=x^{\prime}+\sqrt{ }(2 n+4 \lambda)\left[\frac{1}{8} \frac{c_{3}}{\kappa!}\left(\xi^{2}-1\right)+\right. & \left\{\frac{1}{24} \frac{c_{4}}{\kappa_{2}^{8}}\left(\xi^{3}-3 \xi\right)-\frac{1}{36} \frac{c_{33}}{\kappa_{2}^{3}}\left(2 \xi^{3}-5 \xi\right)\right\}+\left\{\frac{1}{120} \frac{c_{5}}{\kappa!}\left(\xi^{4}-6 \xi^{2}+3\right)\right. \\
& \left.\left.-\frac{1}{24} \frac{c_{34}}{\kappa!}\left(\xi^{4}-5 \xi^{8}+2\right)+\frac{1}{324} \frac{c_{233}}{\kappa!}\left(12 \xi^{4}-53 \xi^{2}+17\right)\right\}\right] . \tag{25}
\end{align*}
$$

In this, $x^{\prime}$ is our first approximation, and the correction improves it considerably even at the lower end of the distribution. The values of the expressions in $\xi$ in (25) are directly available for several probability levels from the table in Cornish \& Fisher's paper.
The approximate values of the probability integral of the $\chi^{\prime \prime}$-distribution obtained by these methods in a few cases are given in Table 4. Table 5 shows the approximate upper and lower $5 \%$ points evaluated by method II.

Comparing the two methods for the probability integral, the second one, employing. terms of the Gram-Charlier series up to $O\left(n^{-1}\right)$, gives greater accuracy and is to be preferred,
although from the point of view of labour and time involved, the first method is simpler and easier to apply. With respect to the percentage points, the method using the Cornish-Fisher inversion appears to be quite good, particularly at the upper points, but it does involve a certain amount of labour.

Table 4. Closer approximations to the $\chi^{\prime 2}$-integral

| $n$ | $\lambda$ | $x$ | 1st approx. | 2nd approx. method |  | Exact |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  | I | III |  |
| 4 | 4 | 10.00 | 0.7191 | 0.7209 | 0.7119 | 0.7118 |
| 4 | 4 | 24.00 | 0.9913 | 0.9817 | 0.9913 | 0.9925 |
| 7 | 16 | 24.00 | 0.5947 | 0.5938 | 0.5889 | 0.6898 |
| 7 | 16 | 38.97 | 0.9482 | 0.9504 | 0.9502 | 0.9500 |
| 16 | 8 | 20.00 | 0.3380 | 0.3345 | 0.3388 | 0.3369 |
| 16 | 8 | 40.00 | 0.9826 | 0.9632 | 0.9631 | 0.9832 |

Table 5. Closer approximation to the $\chi^{\prime 2}$-percentage points, using method II

| $n$ | $\lambda$ | Upper 5\% point |  |  | Lower $5 \%$ point |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | lat approx. | 2nd approx. | Exact | lst approx. | 2nd approx. | Exact |
| 2 | 4 | 14.72 | 14.87 | $14 \cdot 64$ | 0.945 | 0.574 | 0.646 |
| 2 | 18 | 33.35 | 33.06 | 33.06 | 6.891 | 6.526 | 6. 322 |
| 4 | 4 | 17.38 | 17.33 | $17 \cdot 31$ | 1.954 | 1.731 | 1.765 |
| 4 | 16 | 35.60 | 35.42 | 35.43 | $8 \cdot 363$ | 8.017 | 7.884 |
| 7 | 4 | 21.28 | 21-27 | 21-23 | 3.789 | 3.750 | 3-664 |
| 7 | 16 | 39.18 | 38.97 | 38.87 | 10.637 | 10.267 | 10.257 |

## 4. Applioations of the $\chi^{2}$-distribution

## 4-1. The power function of the $\chi^{2}$-test

There are several possible applications of the non-central $\chi^{3}$-distribution in statistics. We shall consider only a few of them. We will show here how this distribution arises in the study of power functions of the $\chi^{3}$-tests and how the approximations of $\S 3$ are useful in this connexiol.

Suppose $\xi_{1}, \xi_{2}, \ldots, \xi_{n}$ are $n$ independent observations in a sample. If we make the null hypothesis $H_{0}$, that the $\xi_{i}$ have been drawn from a normal population with mean zero and s.D. unity, then if $H_{0}$ is true, the statistic $\chi^{2}=\Sigma \xi_{i}^{2}$ will exceed $\chi_{\alpha}^{2}$, the $\alpha$-significance point of the $\chi^{3}$-distribution, based on $n$ degrees of freedom, in a proportion $\alpha$ of the cases.

The power of the $\chi^{3}$-test is given by the probability that $\Sigma \xi^{8}$ exceeds $\chi_{\alpha}^{2}$ under some alternative hypothesis. If as an alternative to $H_{0}$, we suppose that the $\xi_{i}$ have been drawn from normal populations having unit s.d. but different means $a_{i}$, then $\Sigma \xi_{i}^{2}$ will follow the non-
central $\chi^{3}$-distribution with $n$ degrees of freedom and parameter $\lambda=\Sigma a_{i}^{2}$. Denoting this by $p_{n}\left(\chi^{\prime 2} \mid \lambda\right)$, the power function is given by

$$
\begin{equation*}
\int_{\chi_{a}^{\prime}}^{\infty} p_{n}\left(\chi^{\prime 2} \mid \lambda\right) d \chi^{\prime 2} \equiv \beta(n, \lambda, \alpha) . \tag{26}
\end{equation*}
$$

Thus the power is a function of the single parameter $\lambda$ and we may write the null hypothesis as $H_{0}(\lambda=0)$ and an alternative as $H_{1}(\lambda)$, where $H_{1}$ is a composite hypothesis including the family of alternatives for which $\Sigma a_{i}^{q}=\lambda$.

It was shown in $\S 3 \cdot 1$ that the $\chi^{\prime 2}$-distribution is fairly well approximated by a Type III distribution fitted from its first two moments. The power function $\beta$ could therefore be evaluated quickly and fairly accurately by the method of the first approximation. When greater accuracy is needed, one of the other methods described in $\S 3.3$ may be used.

We give here a table (Table B) of values of the power of the $\chi^{2}$-test applied at the significance level $\alpha=0.05$, obtained by the second method of §3.3. The accuracy of these values in different parts of the table can be judged from the closeness between the approximate and exact values of the probability integral shown in Tables 1 and 4. In some of the cases tabled there, the limit $x$ was chosen near to the $5 \%$ point of the corresponding $\chi^{2}$, so as to give a value of

$$
1-\int_{0}^{x} p_{n}(x \mid \lambda) d x
$$

in the neighbourhood of the power $\beta$. It is believed that, in general, there is three-figure accuracy in Table 6.

Table 6. The power function of the $\chi^{2}$-test using a $5 \%$ significance level;
values of $\beta(n, \lambda, \alpha)$, where $\alpha=0.05$

| ${ }^{1} \lambda^{\lambda}$ | 2 | 4 | 6 | 8 | 10 | 12 | 14 | 16 | 18 | 20 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 2 | $0 \cdot 234$ | 0.416 | 0.585 | 0.719 | 0.819 | 0.885 | 0.929 | 0.956 | 0.973 | 0.983 |
| 3 | 0.195 | 0.357 | 0.518 | 0.655 | 0.762 | 0.841 | 0.897 | 0.935 | 0.958 | 0.974 |
| 4 | $0 \cdot 171$ | 0.320 | $0 \cdot 470$ | $0 \cdot 605$ | 0.718 | 0.803 | 0.867 | 0.913 | 0.943 | 0.963 |
| 5 | 0.157 | 0.292 | $0 \cdot 432$ | 0.585 | 0.878 | 0.769 | 0.838 | $0 \cdot 891$ | 0.927 | 0.952 |
| 6 | 0.146 | 0.270 | 0.404 | 0.531 | 0.644 | 0.738 | 0.813 | 0.870 | 0.911 | 0940 |
| 7 | 0.138 | 0.251 | $0 \cdot 378$ | 0.502 | 0.614 | 0.710 | 0.788 | 0.849 | 0.895 | 0.928 |
| 8 | $0 \cdot 131$ | 0.238 | $0 \cdot 357$ | 0.477 | 0.688 | 0.685 | 0.765 | 0.830 | 0.879 | 0.916 |
| $\theta$ | 0.125 | 0.225 | 0.339 | 0.455 | 0.564 | 0.661 | 0.744 | 0.811 | 0.883 | 0.903 |
| 10 | 0.121 | 0.215 | 0.323 | 0.435 | 0.542 | 0.640 | 0.724 | 0.793 | 0.848 | 0.891 |
| 12 | 0.113 | 0.198 | $0 \cdot 297$ | $0 \cdot 402$ | 0.505 | 0.601 | $0 \cdot 688$ | 0.759 | 0.818 | 0.866 |
| 14 | $0 \cdot 108$ | 0.185 | 0.278 | 0.374 | 0.473 | 0.567 | 0.853 | 0.728 | 0.791 | 0.842 |
| 16 | 0.103 | 0.174 | 0.259 | 0.351 | 0.448 | 0.538 | 0.623 | 0.699 | 0.784 | 0.819 |
| 18 | 0.099 | 0.165 | 0.244 | 0.332 | 0.422 | 0.512 | 0.596 | 0.673 | 0.740 | 0.796 |
| 20 | 0.098 | 0.158 | 0.232 | 0.315 | 0.402 | 0.489 | 0.572 | 0.648 | 0.716 | 0.775 |

When $n$ or $\lambda$ is so large that $\nu=n+\lambda^{2} /(n+2 \lambda)$ is over 30 , we may use the normal approximation of $\S 3.2$ for obtaining the power function more quickly than by the method of the $\chi^{2}$-approximation.

The above table can be used in a variety of ways: (a) For given $\lambda$ and $n$, we may ask what is the chance of establishing significance at the $5 \%$ level? ( $b$ ) For given $n$, we may ask how large $\lambda$ must be to have, say, a $90 \%$ chance ( $\beta=0.90$ ) of establishing significance at the $5 \%$ level
when a real difference in the $a_{i}$ exists? (c) For given $\lambda$, we may ask how many observations are necessary to have a chance $\beta$ of establishing significance?

An alternative graphical approach to the inverse problems (b) and (c) is indicated in §7.3, p. 228 below.

### 4.2. Application to the $\chi^{2}$-test for the goodness of fit

The $\chi^{8}$-test for goodness of fit is concerned with the comparison of observed frequencies with those expected under a given hypothesis. The latter may be the theoretical frequencies of a continuous distribution or may be obtained by taking integrals of a continuous frequency distribution over a set of class intervals. Denote the observed frequencies by $n_{i}$ and the expected frequencies by $N \pi_{i}(i=1,2, \ldots, k)$, where $k$ is the number of groups and $N$ the total number of observations in the sample. Then

$$
\begin{equation*}
\sum_{i=1}^{k} n_{i}=\sum_{i=1}^{k} N \pi_{i}=N \tag{27}
\end{equation*}
$$

As is well known, the distribution of

$$
\begin{equation*}
\phi^{2}=\sum_{i=1}^{k} \frac{\left(n_{i}-N \pi_{i}\right)^{2}}{N \pi_{i}} \tag{28}
\end{equation*}
$$

when the $N \pi_{i}$ are the true population expectations, may be related as an approximation to that of the sum of squares of normal variables. To link up also with the non-central theory disoussed in $\S \S 2 \cdot 1-2 \cdot 3$, the following approach may be adopted, although it must be realized that the conclusions reached are not exact. As in all problems concerning $\phi^{2}$, it is generally only possible to assess the degree of error involved, in samples of finite size, by specific numerical comparisons.

As shown originally by K. Pearson ( 1900,1916 ), the variances and co-variances of the $k$ frequencies $n_{i}$, restricted by the condition (27), are precisely those holding in the section

$$
\begin{equation*}
X_{1}+X_{2}+\ldots+X_{k}=0 \tag{29}
\end{equation*}
$$

of the $k$-dimensioned normal probability distribution whose probability density at
is

$$
\begin{gather*}
\left(X_{1}, X_{2}, \ldots, X_{k}\right) \\
p\left(X_{1}, X_{2}, \ldots, X_{n}\right)=\text { constant } \times \exp \left[-\frac{1}{2} \sum_{i} \frac{X_{1}}{N \pi_{i}}\right] . \tag{30}
\end{gather*}
$$

Thus, provided that the expectations $N \pi_{i}$ are large enough to prevent serious inaccuracy from discontinuity effecte or boundary limitations, relationships between the $n_{i}$ may be treated as relationships, within the prime (29), between normal variables $X_{i}$ which in the $k$-dimensioned space are distributed independently with zero means and variances $N \pi_{i}$. With these limitations, we may write

$$
\begin{equation*}
x_{i}=\frac{X_{i}}{\sqrt{\left(N \pi_{i}\right)}}=\frac{n_{i}-N \pi_{i}}{\sqrt{\left(N \pi_{i}\right)}} \quad(i=1, \ldots, k) . \tag{31}
\end{equation*}
$$

The distribution of the $\phi^{2}$ defined in (28) can then be derived from the resulte given in §2.3. The condition $\Sigma n_{i}=N$ may be written

$$
\begin{equation*}
\sum_{i} \sqrt{ } \pi_{i} \frac{n_{i}-N \pi_{i}}{\sqrt{\left(N \pi_{i}\right)}}=0 \tag{32}
\end{equation*}
$$

corresponding to $\sum_{i} c_{i} x_{i}=\rho=0$, where $\sum_{i} c_{i}^{2}=1$. Hence $\phi^{2}$ will be approximately distributed as $\chi^{3}$ with $k-1$ degrees of freedom.

Having in mind the question of the power of the test, we may next ask what will be the distribution of $\phi^{2}$ if the frequencies $N \pi_{i}$ inserted into the expression (28) are not the true expectations? Suppose that $N p_{i}$ are the true expectations; both $\sum_{i} p_{i}$ and $\sum_{i} \pi_{i}$ will be unity.

In the notation of § 2 we now have
while

$$
\begin{gather*}
\xi_{i}=\frac{n_{i}-N \pi_{i}}{\sqrt{\left(N p_{i}\right)}}, \quad x_{i}=\frac{n_{i}-N p_{i}}{\sqrt{\left(N p_{i}\right)}}, \quad a_{i}=\frac{N\left(p_{i}-\pi_{i}\right)}{\sqrt{\left(N p_{i}\right)}},  \tag{33}\\
\sum_{i} \sqrt{ } p_{i} \frac{\left(n_{i}-N \pi_{i}\right)}{\sqrt{ }\left(N p_{i}\right)}=\sum_{i} c_{i} \xi_{i}=0 .
\end{gather*}
$$

It follows that approximately

$$
\begin{equation*}
\phi^{\prime 2}=\sum_{i} \xi_{i}^{2}=\sum_{i} \frac{\left(n_{i}-N \pi_{i}\right)^{2}}{N p_{i}} \tag{35}
\end{equation*}
$$

will be distributed as a non-central $\chi^{2}$ with $k-1$ degrees of freedom and

$$
\begin{equation*}
\lambda^{\prime}=\sum_{i}\left(a_{i}^{2}\right)=N \sum_{i} \frac{\left(p_{i}-\pi_{i}\right)^{2}}{p_{i}} . \tag{36}
\end{equation*}
$$

The sum of squares we need is the $\phi^{2}$ of (28), not the $\phi^{2}$ of (35). By introducing a further approximation we may, however, conclude that $\phi^{2}=\sum_{i}\left(n_{i}-N \pi_{i}\right)^{2} / N \pi_{i}$ is distributed as non-central $\chi^{2}$ with $k-1$ degrees of freedom, and

$$
\begin{equation*}
\lambda=N \sum_{i} \frac{\left(p_{i}-\pi_{i}\right)^{2}}{\pi_{i}} . \tag{37}
\end{equation*}
$$

The approximation involved should not be serious if the differences $\delta_{i}=N \pi_{i}-N p_{i}$ are small compared to $N \pi_{i}$; for

$$
\begin{aligned}
\phi^{\prime 2}=\sum_{i} \frac{\left(n_{i}-N \pi_{i}\right)^{2}}{N p_{i}} & =\sum_{i} \frac{\left(n_{i}-N \pi_{i}\right)^{2}}{N \pi_{i}}\left\{1-\frac{\delta_{i}}{N \pi_{i}}\right\}^{-1} \\
& =\phi^{2}+\sum_{i} \delta_{i} \frac{\left(n_{i}-N \pi_{i}\right)^{2}}{\left(N \pi_{i}\right)^{2}}+\sum_{i} \delta_{i}^{2} \frac{\left(n_{i}-N \pi_{i}\right)^{2}}{\left(N \pi_{i}\right)^{3}}+\ldots
\end{aligned}
$$

Since the multipliers $\delta_{i}$ in the second term may be positive or negative and $\Sigma \delta_{i}=0$, this term will generally be small; the further terms, containing successive powers of $\delta_{i} /\left(N \pi_{i}\right)$, will also be of diminishing importance.

This result makes it possible to determine the power of the goodness of fit test of any simple (completely specified) hypothesis $H_{0}$ (specifying probabilities $\pi_{i}$ ) with respect to a simple alternative hypothesis $H_{1}$ (specifying probabilities $p_{i}$ ). Hence, for any given class of alternatives $H$, we can determine the power function. In so far as the $5 \%$ significance level is used, the power may be determined from Table 6, p. 214, using the $\lambda$ of equation (37) and degrees of freedom $k-1$. Otherwise, we can use the $\chi^{2}$-approximation to the $\chi^{\prime 8}$-distribution developed in §3•1. Thus the power is
where

$$
\begin{gather*}
\int_{\chi_{\alpha}^{2}}^{\infty} p_{k-1}\left(\chi^{\prime 2} \mid \lambda\right) d \chi^{\prime 2}=\int_{\chi_{\alpha}^{\prime} / \rho}^{\infty} p_{v}\left(\chi^{8}\right) d \chi^{2},  \tag{38}\\
\rho=\frac{k-1+2 \lambda}{k-1+\lambda}, \quad \nu=\frac{(k-1+\lambda)^{2}}{k-1+2 \lambda}, \quad \lambda=N\left(\sum_{i} \frac{p_{i}^{2}}{\pi_{i}}-1\right) . \tag{39}
\end{gather*}
$$

[^1]
## P. B. Pathatk

For comparison of this approximate distribution with the exact one, we proceed now to find the exact moments of $\phi^{2}$. It is known (e.g. Haldane, 1937) that under $H_{1}$ the expectations of the powers of the observed frequency $n_{i}$ are
where

$$
\begin{align*}
& \mathscr{E}\left(n_{i}\right)=N p_{i}, \\
& \mathscr{E}(n \boldsymbol{q})=N_{2} p_{i}+N p_{i}, \\
& \mathcal{E}\left(n_{i}^{4}\right)=N_{4} p_{i}^{4}+6 N_{3} p_{i}^{3}+7 N_{2} p_{i}^{\mathbf{4}}+N p^{2}  \tag{40}\\
& \begin{aligned}
& \mathcal{S}\left(n_{i}^{2} n_{j}^{2}\right)=N_{\downarrow} p_{i}^{\frac{2}{2}} p_{j}^{2}+N_{3}\left(p_{i}^{2} p_{j}+p_{j}^{\frac{2}{2}} p_{i}\right)+N_{\mathbf{s}} p_{i} p_{j} \\
& \text { etc. },
\end{aligned} \\
& N_{r}=N!/(N-r)!.
\end{align*}
$$

Writing $\phi^{2}$ in (28) in the form $\quad \phi^{2}=\frac{1}{N} \Sigma\left(n_{i}^{2} / \pi_{i}\right)-N$,
we have

$$
\begin{aligned}
\mathscr{E}\left(\phi^{2}\right) & =\frac{1}{N} \Sigma \frac{\mathscr{E}\left(n_{i}^{2}\right)}{\pi_{i}}-N \\
& =\frac{1}{N}\left\{\Sigma \frac{N_{2} p_{i}^{2}}{\pi_{i}}+\Sigma \frac{N p_{i}}{\pi_{i}}\right\}-N .
\end{aligned}
$$

Hence

$$
\begin{equation*}
\mu_{1}^{\prime}=(N-1) \Sigma\left(p_{i}^{2} / \pi_{i}\right)+\Sigma\left(p_{i} / \pi_{i}\right)-N . \tag{41}
\end{equation*}
$$

Again,

$$
\mathscr{E}\left[\left(\phi^{2}\right)^{2}\right]=\mathscr{E}\left[\frac{1}{N} \Sigma \frac{n_{i}^{4}}{\pi_{i}^{2}}+\frac{1}{N^{2}} \sum_{+j} \frac{n_{i}^{2} n_{j}^{2}}{\pi_{i} \pi_{j}}-2 \Sigma \frac{n_{i}^{2}}{\pi_{i}}+N^{2}\right],
$$

from which on substitution and simplication we obtain

$$
\begin{align*}
& \mu_{2}=N^{-1}\left\{(N-1)(6-4 N)\left[\Sigma\left(p_{i}^{2} / \pi_{i}\right)\right]^{2}+4(N-1)(N-2) \Sigma\left(p_{i}^{3} / \pi_{i}^{2}\right)\right. \\
&-4(N-1) \Sigma\left(p_{i}^{2} / \pi_{i}\right) \Sigma\left(p_{i} / \pi_{i}\right)+6(N-1) \Sigma\left(p_{i}^{2} / \pi_{i}^{2}\right) \\
&\left.-\left[\Sigma\left(p_{i} / \pi_{i}\right)\right]^{2}+\Sigma\left(p_{i} / \pi_{i}^{2}\right)\right\} . \tag{42}
\end{align*}
$$

In a similar way the third moment has also' been obtained but the expression is so long and so difficult to evaluate numerically that it may not be of much value for comparison purposes.

When $p_{i}=\pi_{i}$ the above expressions reduce to those derived by Haldane (1937) for the exact moments of the distribution of $\phi^{2}$ under the null hypothesis.
The approximation to the distribution of $\phi^{2}$ obtained, using the simplification of $\S 3 \cdot 1$, will have the following first two moments:

$$
\left.\begin{array}{l}
\mu_{1}^{\prime}=\nu+\lambda=k-1+\lambda=k-1+N\left[\Sigma\left(p_{i}^{2} / \pi_{i}\right)-1\right]  \tag{43}\\
\mu_{2}=2(\nu+2 \lambda)=2(k-1)+4 \lambda=2(k-1)+4 N\left[\Sigma\left(p_{i}^{2} / \pi_{i}\right)-1\right]
\end{array}\right\}
$$

using the expression for $\lambda$ in (37).
A comparison of these approximate moments with the exact ones, (41) and (42), appears to be only possible numerically. Some comparisons have been made, including a check-up on the whole distribution by a random sampling experiment. In the cases taken, the approximation appeared satisfactory for practical purposes but some further investigation is in hand. The results will be published in a subsequent paper.

### 4.3. Uses of the power function of the $\chi^{2}$ goodness of fit test

We have seen in $\S 4 \cdot 2$ that, to the approximation involved, the power of the $\chi^{2}$-test for $H_{0}$ with regard to an alternative $H_{1}$ is a function of $k-1, \lambda, \alpha$ and can be written $\beta(k-1, \lambda, \alpha)$, where $k$ is the number of groups, $\alpha$ the significance level at which the test is applied and

$$
\lambda=N\left(\sum_{i=1}^{k} \frac{p_{i}^{2}}{\pi_{i}}-1\right)=N \Delta\left(H_{0}, H_{1}\right) .
$$

This shows that $\lambda$ is a function of $\pi_{i}$ and $p_{i}$, and can be regarded as a measure of 'discrepancy' between the two distribution functions specified by $H_{0}$ and $H_{1}$.

The power function can be used to answer several questions connected with the test of goodness of fit: ( $a$ ) For given sample size $N$ and number of groups $k$, we may ask what is the chance of establishing the inadequacy of the hypothesis $H_{0}$, using a given significance level? (b) For given $k$, we may ask how many observations are necessary to give a chance of, say, $90 \%$ of eatablishing significance at the $5 \%$ level? (c) For given $k$ and $N$, we may ask how large a departure of $H_{1}$ from $H_{0}$ (measured by $\Delta\left(H_{0}, H_{1}\right)$ ) will be detected with a given chance?

We shall illustrate these applications by an example from genetics. Consider the intercross

$$
\frac{\mathbf{A B}}{\mathbf{a b}} \times \frac{\mathbf{A B}}{\mathbf{a b}}
$$

where $\mathbf{A}$ and $\mathbf{B}$ are two independent factors, the recessive genes of which are represented by $\mathbf{a}$ and $\mathbf{b}$. The offspring are of the four types [AB], $[\mathbf{A b}],[\mathbf{a B}],[a b]$ with frequencies in the proportions $9,3,3,1$. We test whether the experiment is to confirm this theory or to reject it in favour of a definite alternative giving frequencies proportional to $9,3,3 r, r(r$ being less than 1). This happens when the two classes of offspring containing the two recessive genes ( $a, a$ ) are less viable than those containing only one dominant gene, so that only a fraction of the offspring survive.

Here, the expected frequencies are

Hence

$$
\begin{align*}
& \pi_{i}: 9 / 16,3 / 16,3 / 16,1 / 16 \\
& p_{i}: 9 / 4(3+r), 3 / 4(3+r), 3 r / 4(3+r), r / 4(3+r) \\
& \lambda=N\left(\frac{4\left(3+r^{8}\right)}{(3+r)^{2}}-1\right) \tag{44}
\end{align*}
$$

where $N$ is the number of offspring studied. Then

$$
\begin{equation*}
\Delta \equiv \Delta\left(H_{0}, H_{1}\right)=\frac{4\left(3+r^{2}\right)}{(3+r)^{2}}-1 . \tag{45}
\end{equation*}
$$

Let us now consider the three situations where the power-function idea could be applied.
(a) Suppose we have 100 observations. Using the $\chi^{2}$-test at the $5 \%$ level to test the null hypothesis $(r=1)$, the chance of establishing differential viability when $r=\frac{1}{2}$ is obtained by evaluating $\lambda$ from (37) and then entering Table 6 (p. 214) with this $\lambda$ and $n=k-1=3$. Here $\lambda=300 / 49$ and so the power $\beta=0.52$.
(b) Suppose we want a $90 \%$ chance of detecting that $r=\frac{1}{2}$, using the $5 \%$ significance level. We find from Table 6 that $\lambda=14 \cdot 1$ and hence, putting $r=\frac{1}{2}$ in (45), obtain $\Delta=3 / 49$. Then from (44) we find that we shall need a sample of $N=230$.
(c) Again, if $N=100, \alpha=0.05$, we may ask how small $r$ must be to give a 50 : 50 chance for establishing significance? We find $\lambda$ as before and solve(44)for $r$. Thus taking $\beta=0 \cdot 50$, then $\lambda=5.8$ and $r=0.51$.

### 4.4. A closer approximation to the power function of the $\chi^{2}$ goodness of fit test

In §4.2, when deriving the $\chi^{\prime 8}$-approximation to the distribution of

$$
\phi^{2}=\Sigma \frac{\left(n_{i}-N \pi_{i}\right)^{2}}{N \pi_{i}},
$$

we made the assumption that $\pi_{i}$ and $p_{i}$, the proportions of the expected frequencies under the hypotheses $H_{0}$ and $H_{1}$ do not differ very much, so that we could regard $\left(n_{i}-N p_{i}\right) / \sqrt{ }\left(N \pi_{i}\right)$ as
a normal deviate with zero mean and unit variance. We will now consider the distribution of $\phi^{2}$ without making such an assumption and use it for obtaining a better approximation to the power function.

We can write $\phi^{2}$ in the form

$$
\begin{equation*}
\phi^{2}=\Sigma \frac{p_{i}}{\pi_{i}}\left(\frac{n_{i}-N p_{i}}{\sqrt{\left(N p_{i}\right)}}+\frac{N p_{i}-N \pi_{i}}{\sqrt{\left(N p_{i}\right)}}\right)^{2} \tag{46}
\end{equation*}
$$

the summation being from $i=1$ to $k$. Now, under $H_{1}$ the quantities $\left(n_{i}-N p_{i}\right) / \sqrt{ }\left(N p_{i}\right)$ are distributed approximately normally, as $N(0,1)$, subject to the constraint $\Sigma n_{i}=N$. Hence $\phi^{2}$ in (46) can be regarded as the weighted sum of $k$ normal deviates having different expectations and satisfying the condition $\Sigma n_{i}=N$.

We have obtained in the Appendix (pp. 231-2 below) the characteristic function of the distribution of such a statistic, viz. $\Sigma v_{j}\left(x_{j}+a_{j}\right)^{2}$ subject to the condition $\Sigma c_{j}\left(x_{j}+a_{j}\right)=\rho$. Making the appropriate substitution in (6) of the Appendix, we have the characteristio function of $\phi^{\mathbf{2}}$ :

$$
\begin{align*}
\left(\Sigma \frac{p}{1-2 i t p / \pi}\right)^{-1} & \prod_{1}^{k}(1-2 i t p / \pi)^{-1} \\
& \times \exp \left\{N \Sigma\left(\frac{i t(p-\pi)^{2} / \pi}{1-2 i t p / \pi}\right)-\frac{N}{2}\left(\Sigma \frac{p-\pi}{1-2 i t p / \pi}\right)^{2}\left(\Sigma \frac{p}{1-2 i t p / \pi}\right)^{-1}\right\} \tag{47}
\end{align*}
$$

where the subscripts of $p_{i}$ and $\pi_{i}$ are dropped. From this the expressions for the first three moments are derived. Thus

$$
\left.\begin{array}{rl}
\mu_{1}^{\prime}=(N-1) \Sigma\left(p_{i}^{2} / \pi_{i}\right)+\Sigma\left(p_{i} / \pi_{i}\right)-N \\
\mu_{2}=4(N-1) \Sigma\left(p_{i}^{8} / \pi_{i}^{q}\right)-2(2 N-1)\left[\Sigma\left(p_{i}^{\mathbf{2}} / \pi_{i}\right)\right]^{2}+2 \Sigma\left(p_{i}^{2} / \pi_{i}^{q}\right), \\
\mu_{3}=24(N-1) \Sigma\left(p_{i}^{4} / \pi_{i}^{\mathbf{2}}\right)-24(2 N-1)\left[\Sigma\left(p_{i}^{2} / \pi_{i}\right)\right]\left[\Sigma\left(p_{i}^{3} / \pi_{i}^{\mathbf{q}}\right)\right]  \tag{48}\\
& +8(3 N-1)\left[\Sigma\left(p_{i}^{2} / \pi_{i}\right)\right]^{3}+8 \Sigma\left(p_{i}^{\mathbf{3}} / \pi_{i}^{3}\right)
\end{array}\right\}
$$

It will be seen that the only assumption made here, that $\left(n_{i}-N p_{i}\right) / \sqrt{ }\left(N p_{i}\right)$ is distributed as $N(0,1)$ under $H_{1}$, is parallel to the assumption on which the $\chi^{8}$-test of goodness of fit is based, namely, that $\left(n_{i}-N \pi_{i}\right) / \sqrt{ }\left(N \pi_{i}\right)$ is distributed as $N(0,1)$ under $H_{0}$, which is justified when $N \pi_{i}$ are not too small. So, when $N p_{i}$ are not too small we can expect the momenta in (48) to agree well with the true moments (the first two of which are given in (41) and (42)). Obviously the expressions for $\mu_{1}^{\prime}$ are identical. The values of $\mu_{2}$ in the cases examined in the investigation referred to on p. 217 were found to be very close.

We may now obtain a representation of the distribution of $\phi^{2}$ under $H_{1}$ as a Type III having the first two moments of (48), that is, assume $\phi^{2} / \rho$ as distributed as $\chi^{2}$ with $\nu$ degrees of freedom, where $\rho=\frac{1}{2} \mu_{2} / \mu_{1}^{\prime}, \nu=2 \mu_{1}^{\prime 2} / \mu_{2}$. Clearly this will be a better approximation than that of the Type III fitted from the $\mu_{1}^{\prime}, \mu_{2}$ given in (43), and the power function based on this will be closer to the exact one than that based on (38) and (39). But, although there is gain in accuracy, the simplicity of the approximate method is lost. We may similarly consider fitting a Type III distribution, using the true $\mu_{1}^{\prime}$ and $\mu_{2}$, but the labour of computation of $\mu_{2}$, given in (42), appears to be prohibitive.

## 5. CONDITIONAL POWRE FUNOTIONS

In § 4 we have considered the power function of the $\chi^{2}$ goodness of fit test when the null hypothesis is fully specified, i.e. is a simple hypothesis. But often we are interested in testing whether an observed sample has come from a certain type of population, so that we are given
only the form of the population law, not the values of its parameters, say $\theta_{1}, \theta_{\mathbf{2}}, \ldots, \theta_{r} . H_{0}$ is then a composite hypothesis. Sometimes, also, we have to test the hypothesis that several samples are from the same population, without specifying anything about it. In these cases we obtain estimates of the unspecified parameters, say $T_{1}, T_{2}, \ldots, T_{r}$, from the sample and hence calculate the expected cell frequencies $\hat{m}_{i}$. Then, if the method of estimation is efficient $\dagger$,

$$
\begin{equation*}
\phi^{2}=\sum_{i}\left(n_{i}-\hat{m}_{i}\right)^{2} / \hat{m}_{i} \tag{49}
\end{equation*}
$$

is known still to follow approximately a $\chi^{2}$-distribution with $k-r-1$ degrees of freedom.
Suppose now that as alternative to the composite hypothesis $H_{0}$, there is a simple hypothesis $H_{1}$. The question then arises: By estimating the $m_{i}$ on the assumption that $H_{0}$ is true and applying the $\chi^{2}$-test, what chance have we of rejecting $H_{0}$, when, in fact, $H_{1}$ is true?

Some consideration has been given to this problem, and it seems possible to obtain a solution by making use, as a first step, of what David (1947, p. 339) has termed the conditional power function. This gives the chance of rejecting $H_{0}$ when the test is confined to a restricted set, $S$, of samples which provide the same values, say $T_{1}^{(s)}, T_{8}^{(8)}, \ldots, T_{r}^{(s)}$ for the estimated parameters. Thus, if the process of fitting involves estimating two parameters from the sample mean and variance, samples of a set would be those having a common mean and variance. Again, in testing for independence in a contingency table, the conditional power function would be obtained for a set of samples giving the same marginal totals (see Patnaik, 1948). The development of this method will be left for a later communication.

## 6. The non-central $\boldsymbol{F}$-distribution and approximations to it

Suppose two independent variates, $\chi_{1}^{\prime 2}$ and $\chi_{3}^{9}$, follow respectively a non-central $\chi^{2}$-distribution with degrees of freedom $\nu_{1}$ and parameter $\lambda$ and a $\chi^{2}$-distribution with degrees of freedom $\nu_{2}$. Then the ratio

$$
F^{\prime}=\frac{\chi_{1}^{\prime 2} / \nu_{1}}{\chi_{2}^{2} / \nu_{2}}
$$

will have the following probability distribution:

$$
\begin{equation*}
p\left(F^{\prime}\right)=\sum_{j=0}^{\infty}\left[\frac{e^{-i \lambda}\left(\frac{1}{2} \lambda\right)^{\prime}}{j!B\left(\frac{1}{2} \nu_{1}+j, \frac{1}{2} \nu_{2}\right)}\left(\frac{\nu_{1}}{\nu_{2}}\right)^{i \nu_{1}+j} \cdot F^{\prime i \nu_{1}-1+j}\left(1+\frac{\nu_{1}}{\nu_{2}} F^{\prime}\right)^{-i\left(\nu_{1}+\nu_{2}\right)-i}\right], \tag{50}
\end{equation*}
$$

which may be termed the distribution of non-central $\boldsymbol{F}$ or of $\boldsymbol{F}^{\prime}$. This corresponds to Fisher's distribution $C$ (1928). Wishart (1932) considered it in the form of the distribution of the correlation ratio

$$
E^{2}=\frac{\nu_{1} F^{\prime}}{\nu_{2}+\nu_{1} F^{\prime}}
$$

Later, Tang (1938) derived the same from that of $\chi^{\prime 2}$.
If in (50) we put $\nu_{1}=1$, then it reduces to the distribution of non-central $t^{2}$. Denoting the non-central $t$ by $t^{\prime}$, we have

$$
t^{\prime}=\frac{z+\delta}{\sqrt{w}}
$$

where $z$ is a normal deviate with expected value zero and $w$ is an unbiased estimate of its variance. Neyman (1935), Neyman \& Tokarska (1936) and Johnson \& Welch (1939) have
$\dagger$ I. e. gives a solution not very different from the maxintum likelihood or minimum $\chi^{\mathbf{2}}$ solutions, which are nearly identical in large asmples.
dealt with this distribution in detail and studied its various applications. We will not therefore consider here in particular this special case of the $F^{\prime \prime}$-distribution.

Taking the general form (50), we may, by analogy, call $\nu_{1} ; \nu_{2}$ the degrees of freedom and $\lambda$ the non-central parameter. It can be seen that when $\nu_{2}$ tends to infinity the distribution of $F^{\prime}$ reduces to that of $\chi_{1}^{\prime 2} / \nu_{1}$.

The characteristic function is obtained as an infinite sum of confluent hypergeometric functions

$$
\begin{aligned}
\phi(t)=e^{-1 \lambda}\left[H\left(\frac{\nu_{1}}{2},-\frac{\nu_{2}}{2},-\frac{\nu_{2}}{\nu_{1}} i t\right)+\frac{\lambda}{2}\right. & H\left(\frac{\nu_{1}}{2}+1,-\frac{\nu_{9}}{2},-\frac{\nu_{2}}{\nu_{1}} i t\right) \\
& \left.+\left(\frac{\lambda}{2}\right)^{2} \frac{1}{2!} H\left(\frac{\nu_{1}}{2}+2,-\frac{\nu_{2}}{2},-\frac{\nu_{2}}{\nu_{1}} i t\right)+\ldots\right]
\end{aligned}
$$

in which the function, $H(a, b, x)$, is the sum of the series

$$
1+\frac{a}{b} x+\frac{a(a+1)}{2!b(b+1)} x^{2}+\ldots
$$

Thence we derive the following expressions for the first four moments about the origin:

$$
\left.\begin{array}{l}
\mu_{1}^{\prime}=\frac{\nu_{2}\left(\nu_{1}+\lambda\right)}{\left(\nu_{2}-2\right) \nu_{1}}, \\
\mu_{2}^{\prime}=\frac{\nu_{2}^{2}}{\left(\nu_{2}-2\right)\left(\nu_{2}-4\right) \nu_{1}^{2}}\left[\left(\nu_{1}+\lambda\right)^{2}+2\left(\nu_{1}+2 \lambda\right)\right], \\
\mu_{3}^{\prime}=\frac{\nu_{8}^{3}}{\left(\nu_{2}-2\right)\left(\nu_{2}-4\right)\left(\nu_{2}-6\right) \nu_{1}^{3}}\left[\left(\nu_{1}+\lambda\right)^{3}+6\left(\nu_{1}+\lambda\right)\left(\nu_{1}+2 \lambda\right)+8\left(\nu_{1}+3 \lambda\right)\right],  \tag{51}\\
\mu_{4}^{\prime}=\frac{\nu_{2}^{4}}{\left(\nu_{2}-2\right)\left(\nu_{2}-4\right)\left(\nu_{2}-6\right)\left(\nu_{2}-8\right) \nu_{1}^{4}}\left[\left(\nu_{1}+\lambda\right)^{4}+12\left(\nu_{1}+\lambda\right)^{2}\left(\nu_{1}+2 \lambda\right)\right. \\
\\
\\
\left.\quad+44\left(\nu_{1}+2 \lambda\right)^{2}+48\left(\nu_{1}+4 \lambda\right)-32 \lambda^{2}\right],
\end{array}\right\}
$$

of which the first two were obtained by Wishart by a different method.
Methods of evaluating the probability integral of the $F^{\prime}$-distribution have been worked out by Wishart and Tang. They involve a considerable amount of labour. Following the procedure adopted in the case of $\chi^{\prime 2}$, it may be possible to obtain a quick, though approximate, method by fitting an $F$-distribution with the exact first two moments of $F^{\prime}$. If we regard $F^{\prime} \mid k$ as following an $F$-distribution with $\nu$ and $\nu_{2}$ degrees of freedom, then, equating the expressions for $\mu_{1}^{\prime}$ and $\mu_{2}$, we have

$$
\begin{gathered}
\frac{\nu_{2}\left(\nu_{1}+\lambda\right)}{k\left(\nu_{2}-2\right) \nu_{1}}=\frac{\nu_{2}}{\nu_{2}-2}, \\
\frac{\nu_{2}^{2}}{k^{2}\left(\nu_{2}-2\right)\left(\nu_{2}-4\right) \nu_{1}^{2}}\left[\left(\nu_{1}+\lambda\right)^{2}+2\left(\nu_{1}+2 \lambda\right)\right]=\frac{\nu_{2}^{2}}{\left(\nu_{2}-2\right)\left(\nu_{2}-4\right) \nu+2},
\end{gathered}
$$

which give the scale factor and the modified degrees of freedom, viz.

$$
\begin{equation*}
k=\frac{\nu_{1}+\lambda}{\nu_{1}}, \quad \nu=\frac{\left(\nu_{1}+\lambda\right)^{2}}{\nu_{1}+2 \lambda} \tag{52}
\end{equation*}
$$

The same result will follow if we approximate the distribution of $\chi_{1}^{\prime 2}$ (the numerator in $F^{\prime}$ ) by a Type III from the first two moments as in §3.1.

Using the above approximation, the probability integral

$$
\int_{0}^{F^{\prime}} p_{r_{1}, r_{2}}\left(F^{\prime} \mid \lambda\right) d F^{\prime}
$$

is approximately equal to

$$
\int_{0}^{F^{\prime} / k} p_{r, p_{\mathbf{2}}}(F) d F
$$

where $k$ and $\nu$ are defined in (52). This can be expressed in the form of an Incomplete $B$-function, viz.
where

$$
\begin{gathered}
I_{x}\left(\frac{\nu}{2}, \frac{\nu_{\mathrm{g}}}{2}\right), \\
x= \\
\frac{\nu F^{\prime} \mid k}{\nu_{2}+\nu F^{\prime} \mid k} .
\end{gathered}
$$

For given values of $\nu_{1}, \nu_{8}, \lambda$ and $F^{\prime}$, we can therefore evaluate the integral from the Tables of the Incomplete B-function (K. Pearson, 1934). When $\nu_{1}$ is even or, if odd, is less than 22, we need interpolate only for $x$ and $\frac{1}{2} \nu(=p)$. Four-point Lagrangian interpolation $p$-wise and linear interpolation $x$-wise will be necessary.

Tang's tables of $P_{\text {II }}$ (the error of the second kind) (1938) give exact values of the integral of the $E^{2}$-distribution, which, put in the $F^{\prime}$-form, is

$$
\begin{equation*}
\int_{0}^{F_{a}} p_{r_{1}, r_{z}}\left(F^{\prime} \mid \lambda\right) d F^{\prime}, \tag{53}
\end{equation*}
$$

Table 7. Approximate and exact values of the $F^{\prime}$-integral, $\int_{0}^{x} p_{\nu_{1}, s_{s}}\left(F^{\prime} \mid \lambda\right) d F^{\prime}$

| $\nu_{1}$ | $\nu_{1}$ | $\lambda$ | $x$ | Approx. | Exact |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 3 | 10 | $\begin{array}{r} 4 \\ 4 \\ 16 \\ 16 \end{array}$ | $\begin{aligned} & \mathbf{3 . 7 0 8} \\ & \mathbf{6} 552 \\ & \mathbf{3 . 7 0 8} \\ & \mathbf{6 . 5 5 2} \end{aligned}$ | $\begin{aligned} & 0.752 \\ & 0.019 \\ & 0.203 \\ & 0.520 \end{aligned}$ | $\begin{aligned} & 0.745 \\ & 0.918 \\ & 0.208 \\ & 0.617 \end{aligned}$ |
| 3 | 20 | $\begin{array}{r} 4 \\ 4 \\ 16 \\ 16 \end{array}$ | $\begin{aligned} & \mathbf{3 . 0 9 8} \\ & \mathbf{4 . 9 3 8} \\ & \mathbf{3 . 0 9 8} \\ & 4.938 \end{aligned}$ | $\begin{aligned} & 0.706 \\ & 0.889 \\ & 0.119 \\ & 0.350 \end{aligned}$ | $\begin{aligned} & 0.700 \\ & 0.887 \\ & 0.126 \\ & 0.347 \end{aligned}$ |
| 5 | 10 | $\begin{array}{r} 6 \\ 6 \\ 24 \\ 24 \end{array}$ | $\begin{aligned} & \mathbf{3 . 3 2 6} \\ & 5.636 \\ & 3.328 \\ & 5.636 \end{aligned}$ | $\begin{aligned} & 0.731 \\ & 0.913 \\ & 0.157 \\ & 0.463 \end{aligned}$ | $\begin{aligned} & 0.731 \\ & 0.914 \\ & 0.158 \\ & 0.461 \end{aligned}$ |
| 5 | 20 | 6 6 24 $\mathbf{2 4}$ | $\begin{aligned} & 2 \cdot 711 \\ & 4 \cdot 103 \\ & 2 \cdot 711 \\ & 4 \cdot 103 \end{aligned}$ | $\begin{aligned} & 0.665 \\ & 0.869 \\ & 0.064 \\ & 0.244 \end{aligned}$ | $\begin{aligned} & 0.664 \\ & 0.870 \\ & 0.069 \\ & 0.245 \end{aligned}$ |
| 8 | 10 | 9 9 38 38 | $\begin{aligned} & 3.072 \\ & 5.057 \\ & 3.072 \\ & 5.057 \end{aligned}$ | $\begin{aligned} & 0.715 \\ & 0.909 \\ & 0.117 \\ & 0.409 \end{aligned}$ | $\begin{aligned} & 0.714 \\ & 0.908 \\ & 0.119 \\ & 0.408 \end{aligned}$ |
| 8 | 30 | 9 8 36 36 | $\begin{aligned} & 2.286 \\ & 3 \cdot 173 \\ & 2.286 \\ & 3 \cdot 173 \end{aligned}$ | $\begin{aligned} & 0.581 \\ & 0.815 \\ & 0.014 \\ & 0.085 \end{aligned}$ | $\begin{aligned} & 0.578 \\ & 0.813 \\ & 0.017 \\ & 0.088 \end{aligned}$ |

$F_{\alpha}$ being the $\alpha$-percentage point of the $\boldsymbol{F}$-distribution with $\nu_{1}, \nu_{2}$ degrees of freedom. Two levels of $\alpha$ were chosen for the tables, namely, 0.05 and 0.01 , and the range of $\nu_{1}$ is 0 to 8 . The tables have to be entered with $\phi=\sqrt{ }\left[\lambda /\left(\nu_{1}+1\right)\right]$. Since $\phi$ is at intervals of $0 \cdot 5$, the corresponding intervals for $\lambda$ are very wide, which therefore makes interpolation unsatisfactory.

Table 7 gives the values of the integral (53) calculated by the approximare method indicated above, for certain cases where Tang's exact values are available. The comparison shows that, in general, we have two-figure accuracy, while the error in the third place appears to be quite small near the tails. $\dagger$

It is to be noted that the table compares the integral at only two points, the 5 and $1 \%$ points of the corresponding $\boldsymbol{F}$-distribution. Due to the lack of exact values it has not been possible to judge the closeness at other points. However, some idea of the general accuracy could be had by comparing the true and approximate figures for different $\lambda$ 's with the same $\nu_{1}, \nu_{2}$ and $x\left(=F_{a}\right)$.

It can be easily shown (see Hartley, 1948) that the maximum error in the $F^{\prime}$-integral due to our approximation will not exceed the maximum error in the corresponding $\chi^{\prime 2}$-integral, that is, in

$$
\int_{0}^{x^{\prime \prime}} p_{v_{1}}\left(\chi^{\prime 2} \mid \lambda\right) d \chi^{\prime 2}
$$

Table 1 on p. 207 gives an idea of the magnitude of the errors in the $\chi^{\prime 2}$-integral, and so we can say that the errors in the $F^{\prime}$-integral will not be of a higher order.
The percentage points of $F^{\prime}$ can be obtained by interpolation in the $F$-tables (Merrington \& Thompson, 1843), for the fractional $\nu$ and $\nu_{2}$ and multiplying the interpolate by $k$ in (52).

Closer approximations to the probability integral and percentage points may be derived by the method based on the Gram-Charlier series, analogous to the second method of $\S 3 \cdot 3$.

## 7. The power funotion of the analysis of variance tests

### 7.1. Evaluation of the power function

The test of a general linear hypothesis may be formulated as follows: Suppose $x_{1}, x_{2}, \ldots, x_{N}$ be $N$ normal variates with means $\xi_{1}, \xi_{2}, \ldots, \xi_{n}$ and the same s.D., $\sigma . \xi_{i}$ is a linear function of $s<N$ parameters, $\theta_{1}, \theta_{2}, \ldots, \theta_{8}$. Thus

$$
\xi_{i}=a_{i 1} \theta_{1}+a_{i 8} \theta_{\mathbf{z}}+\ldots+a_{i s} \theta_{t}
$$

The linear hypothesis specifies, say, $r$ of these parameters, i.e.

$$
\begin{equation*}
\theta_{1}=\theta_{1}^{0}, \quad \theta_{\mathbf{2}}=\theta_{2}^{0}, \quad \ldots, \quad \theta_{r}=\theta_{r}^{0} . \tag{54}
\end{equation*}
$$

It is possible by a suitable transformation of variates (see Tang, 1038) of the form
to transform

$$
y_{j}=c_{j 1} x_{1}+c_{j 2} x_{8}+\ldots+c_{j N} x_{N}
$$

into

$$
T_{i-s+r}^{2}=\sum_{i=1}^{N}\left(x_{i}-\xi_{i}\right)^{2}
$$

$$
T^{2}=\sum_{j=1}^{N-s} y_{j}^{2}+\sum_{j-N-s+1}^{N-s+r}\left(y_{j}-\eta_{j}\right)^{2}+\sum_{j=N-s+r+1}^{N}\left(y_{j}-\eta_{j}\right)^{2},
$$

where $\eta_{j}$ in the second sum is a linear function of $\theta_{1}^{0}, \theta_{8}^{0}, \ldots, \theta_{r}^{0}$ and $\eta_{j}$ in the third sum is a linear function of all the $\theta$ 's, while the $a$ 's and c's enter as coefficients.

[^2]To test the hypothesis (54) we consider the oriterion

$$
\begin{equation*}
\left\{\frac{T_{\min .}^{\mathbf{s}}\left(\theta_{1}^{0}, \ldots, \theta_{r}^{0}, \theta_{r+1}, \ldots, \theta_{s}\right)}{T_{\min .}^{\mathrm{s}} .\left(\theta_{1}, \ldots, \theta_{r}, \ldots, \theta_{s}\right)}-1\right\}=\sum_{j=N-s+1}^{N-s+r}\left(y_{j}-\eta_{j}\right)^{2} / \sum_{j=1}^{N-s} y_{j}^{s} . \tag{55}
\end{equation*}
$$

If the hypothesis specifies such values for $\theta_{1}, \theta_{2}, \ldots, \theta_{r}$ that $\eta_{j}$ 's in (55) vanish, then the numerator and denominator are the sums of $r$ and $N-s$ central squares respectively. So, the ratio of the mean squares follows an $F$-distribution. On the other hand, if the $\eta_{j}$ 's do not all vanish, we have the ratio of a sum of $r$ non-central squares to the sum of $N-s$ central squares; hence, the ratio of the mean squares is distributed as non-central $F$, the parameter $\lambda$ being $\Sigma \eta_{j}^{2}$ which can be expressed in terms of $\theta_{1}^{o}, \ldots, \theta_{r}^{0}$ (see Tang, p. 137). Thus we get the $F$-test of the analysis of variance and obtain the power function of this test with respect to an alternative hypothesis as an $F^{\prime}$-integral.

We shall now consider the question of evaluating the power of the analysis of variance test by taking as an illustration the simple case of $k$ groups of observations

$$
\begin{gather*}
x_{t i}(i=1, \ldots, n ; t=1, \ldots, k) . \\
x_{t i}=A+B_{t}+z_{t i}, \tag{56}
\end{gather*}
$$

where $A$ is the general mean, $B_{i}$ the deviation of the mean of the th group from the general mean so that $\Sigma B_{1}=0$ and $z_{t i}$ 's are random residuals, distributed normally with mean zero and s.d. $=\sigma_{0}$. The expressions for the mean squares between groups and within groups follow from the set-up (56):

$$
\begin{aligned}
v & =\frac{1}{k-1} \sum_{t=1}^{k}\left(\bar{x}_{t} .-\bar{x} . .\right)^{2}=\frac{1}{k-1} \sum_{t=1}^{k} n\left(\bar{z}_{t} .-\bar{z} . .+B_{t}\right)^{\mathbf{z}}, \\
v_{0} & =\frac{1}{k(n-1)} \sum_{t=1}^{k} \sum_{i=1}^{n}\left(x_{i t}-\bar{x}_{t} \cdot\right)^{2}=\frac{1}{k(n-1)} \sum_{t=1}^{k} \sum_{i=1}^{n}\left(z_{z_{i}}-\bar{z}_{t} .\right)^{2},
\end{aligned}
$$

where the symbols have the usual meanings. Since ( $\bar{z}_{4},-\bar{z}$..) is a normal deviate with zero mean and variance $\sigma_{0}^{2} / n$, we see that $v$ is the sum of $k$ non-central squares subject to the linear constraint

$$
\sum_{t=1}^{k}\left(\bar{z}_{t}--\bar{z}_{. .}+B_{t}\right)=0
$$

Since further $\Sigma B_{i}=0$, we find from the result of $\S 2 \cdot 3$ that $v$ is distributed as $\sigma_{0}^{2} \chi_{1}^{\prime 2} /(k-1)$, where $\chi_{1}^{\prime 2}$ has $(k-1)$ degrees of freedom and parameter

## Writing

$$
\begin{align*}
& \lambda=n \Sigma B_{i}^{9} / \sigma_{0}^{\mathbf{8}} . \\
& S^{\mathbf{z}}=\left(\Sigma B_{i}^{2}\right) / k \tag{57}
\end{align*}
$$

for the variability between the groups, we have

$$
\begin{equation*}
\lambda=k n S^{2} / \sigma_{0}^{2} . \tag{58}
\end{equation*}
$$

Now $v_{0}$ follows the distribution of $\sigma_{0}^{8} \chi_{\mathbf{3}}^{\mathbf{3}} /[k(n-1)]$, where $\chi_{2}^{\mathbf{3}}$ has $k(n-1)$ degrees of freedom. Hence $v / v_{0}$ is distributed as

$$
\frac{1}{k-1} \chi_{1}^{\prime 2} / \frac{1}{k(n-1)} \chi_{2}^{2},
$$

i.e. as $F^{\prime}$ with $\nu_{1}=k-1, \nu_{2}=k(n-1)$ and $\lambda$ given by (58).

In this example we desire to test for any possible difference between the averages of the groups, so that our null hypothesis is

$$
\begin{equation*}
B_{1}=B_{s}=\ldots=B_{k}=0 . \tag{59}
\end{equation*}
$$

Then, from (57), $S^{\mathbf{2}}$ and therefore $\lambda$ is zero. Hence $v / v_{0}$ follows an $F$-distribution and we get an $F$-test. Thus the test of the hypothesis in (59) is based on the critical region

$$
\begin{equation*}
\frac{v}{v_{0}} \geqslant F_{a}, \tag{60}
\end{equation*}
$$

where $\alpha$ is the significance level at which we are tesiting.
Let us consider an alternative hypothesis that the $B_{i}$ 's are not all zero. Then it is known that the power function, that is, the probability that $\left(v / v_{0}\right) \geqslant F_{\alpha}$, depends only on the single parameter

$$
\frac{S^{2}}{\sigma_{0}^{2}}=\frac{\Sigma B_{i}^{2}}{k \sigma_{0}^{2}} .
$$

Hsu (1941) has shown that amongst all critical regions of $\operatorname{size} \alpha$, whose power functions depend on the single parameter ( $S^{2} / \sigma_{0}^{2}$ ), the critical region of ( 60 ) is the most powerful.

Thus we specify the hypothesis alternative to the null hypothesis (59) by the single parameter $S^{2} / \sigma_{0}^{2}$ in place of the individual parameters, the $B_{i}$ 's. In certain situations, as, for instance, in a manufacturing process, we are more interested in detecting the over-all variability in a set of machines than in detecting the deviation of each particular machine from the general machine average. Then the power function will be useful in measuring the chance of detecting this over-all variability by means of the $F$-test.

The power function of the analysis of variance tests has been considered by Tang (1938) and Hsu (1941). The rather restricted scope of Tang's tables has already been mentioned in $\S 6$. The labour involved in computing the exact values of the power is very heavy, and no tabling on an extensive scale has so far been found possible. However, with the approximations to the $F^{\prime}$-distribution derived in $\S 6$, we may obtain easily a sufficiently accurate value for the power function of the test of any linear hypothesis.
Returning to the case of $k$ groups and $k n$ observations, we have the power function given by

$$
\beta\left(\frac{S^{2}}{\sigma_{0}^{2}}\right)=\int_{F_{a}}^{\infty} p_{v_{1}, v_{\mathbf{a}}}\left(F^{\prime} \mid \lambda\right) d F^{\prime \prime}
$$

where $F_{\alpha}$ is the $\alpha$ percentage point of the $F$-distribution with degrees of freedom $\nu_{1}, \nu_{\mathbf{2}}$. Following the procedure of $\S 6$, this integral approximately equals
where

$$
\begin{gather*}
\int_{F_{a} v_{1}\left(v_{1}+\lambda\right)}^{\infty} p_{r, p_{2}}(F) d F,  \tag{61}\\
\nu=\frac{\left(\nu_{1}+\lambda\right)^{2}}{\nu_{1}+2 \lambda} .
\end{gather*}
$$

Therefore, to this approximation, we have
in which

$$
\left.\begin{array}{l}
\beta\left(\frac{S^{2}}{\sigma_{0}^{2}}\right)=I_{x}\left(\frac{1}{\frac{1}{2}} \nu_{2}, \frac{1}{2} \nu\right)  \tag{62}\\
x=\frac{\left(\nu_{1}+2 \lambda\right) \nu_{2}}{\left(\nu_{1}+2 \lambda\right) \nu_{2}+\left(\nu_{1}+\lambda\right) \nu_{1} F_{a}}
\end{array}\right\}
$$

### 7.2. The difference between systematic and random effects

Next we shall consider two alternatives that arise in practical situations-the random and systematic set-ups (see Daniels, 1939) which may best be described in terms of two examples:

If the groups in the previous illustration correspond to villages and the observations are the
yields of fields in a crop survey, then we can regard the $k$ villages as a random sample from a population of villages and the random set-up represented by

$$
\begin{equation*}
x_{t i}=A+y_{t}+z_{i i} \tag{83}
\end{equation*}
$$

becomes relevant. Here, $A$ is the general mean, $y_{t}$ 's are the group means which are independent random variables with expected value zero and s.D. $=\sigma$, and $z_{4 t}$ 's the random residuals having mean zero and s.d. $=\sigma_{0}$.
On the other hand, if the groups correspond to $k$ machines whioh, from the user's standpoint, constitute the entire population of machines, we cannot regard them as a sample, and so the systematic set-up, in (56), considered on p. 224, is relevant. The null hypothesis in the random set-up is that the parameter $\sigma^{8}=0$, and in the other that $S^{2}=0$ (which is equivalent to (59)). But it is easily seen that both lead to the same $F$-test for the null hypothesis.
In applying the test, we are on the look out for the existence of alternative conditions, where in one case $\sigma^{8}$ and in the other $S^{2}$ is $>0$. It will be noted that ( $S^{2} / \sigma_{0}^{2}$ ) of the systematic set-up corresponds to ( $\sigma^{2} / \sigma_{0}^{2}$ ) of th.c random set-up. Both are measures of relative variability between groups and may be termed 'relative group variability'.
It is possible to relate the power function under the random set-up to that under the systematic set-up. If we regard the $k$ groups as a sample from an infinite number of groups, then $\Sigma B_{l}^{2} /(k-1)$, i.e. $k S^{2} /(k-1)$ will be the sample estimate of the population variance $\sigma^{2}$. Thus treating $S^{2}$ as a random variable having a probability distribution denoted by $p\left(S^{2} / \sigma^{2}\right)$, we can obtain the average power over all the $S^{2}$ s. Thus

$$
\beta=\int_{0}^{\infty} \beta\left(\frac{S^{2}}{\sigma_{0}^{2}}\right) p\left(S^{2} \mid \sigma^{2}\right) d S^{2}
$$

gives the power when the random set-up applies.
This power $\beta$ for given ( $\sigma^{2} / \sigma_{0}^{2}$ ) is directly obtained (see Johnson, 1948) from the $F$-integral:

$$
\begin{equation*}
\int_{F_{a} /\left(n \sigma^{2} / \sigma_{0}^{2}+1\right)}^{\infty} p_{p_{1}, v_{1}}(F) d F=\int_{F_{a}\left(\nu_{1}+1\right) /\left(\nu_{1}+1+\lambda\right)^{\infty}}^{\infty} p_{v_{1}, v_{2}}(F) d F, \tag{64}
\end{equation*}
$$

where

$$
\nu_{1}=k-1, \quad \nu_{8}=k(n-1) \quad \text { and } \quad \lambda=k n \sigma^{3} / \sigma_{0}^{2} .
$$

This can be put in the form of the Incomplete B-function
where

$$
\begin{equation*}
\left.x^{\prime}=\frac{I_{x^{\prime}\left(\frac{1}{2} \nu_{3}, \frac{1}{2} \nu_{1}\right),}}{\left(\nu_{1}+1+\lambda\right) \nu_{2}+\left(\nu_{1}+1\right) \nu_{3} F_{\alpha}} .\right\} \tag{65}
\end{equation*}
$$

It is interesting to note a result which we believe is true in general and whioh on intuitional grounds might be expected to hold, namely, if the null hypothesis is not true, then for the same numerical values of the ratios $S^{2} / \sigma_{0}^{2}$ and $\sigma^{2} / \sigma_{0}^{2}$, the power of the $F$-test is greater in the systematic case than in the random. Four particular cases have been examined numerically as follows:

|  | (a) | (b) | (c) | (d) |
| :---: | :---: | :---: | :---: | :---: |
| Number of groups,,$k$ <br> Number of observations <br> in each group, $n$ | 4 | 11 | 12 | 10 |

Values of the power have been calculated, using equations (62) and (65), and are plotted in Fig. $2(a)-(d)$ as ordinates against $\sigma^{2} / \sigma_{0}^{8}\left(=S^{2} / \sigma_{0}^{3}\right)$. We find from these that the systematic power curve lies above the other; further, we note that the curves are closer to one another in (c) and (d) than in (a) and (b), a fact which agrees with theory that the two power functions must tend to each other with increasing $k$. The errors of approximation in calculating the power in the systematic case are likely to be small judged by the comparative Table 7 and should not affect the relative positions of the power curves.


Fig. 2. Power curves for the random and systematic set-ups for $k$ groups with $n$ observations in each: - - - random, _-_ systematic.

This relation may be interpreted in a different way. Taking case (a) above, it will be seen from Fig. $2(a)$ that we can detect, for instance, a 'systematic' relative group variability of 0.45 with a $70 \%$ chance, while we cannot, with the same chance, detect a variability of magnitude less than 0.9 in the random case. The difference is of course to be expected. For the random set-up, our appreciation of $\sigma^{2}$ is obscured by random variations in both $y$ and $z$ of equation (63); for the systematic set-up, our appreciation of $S^{\mathbf{2}}$ is only obscured by random fluctuations in the $z$ of equation (56).

### 7.3. Applications of the power-function

We will be concerned here mainly with the systematic set-up and will illustrate the application of our results, taking the simple case of $k$ groups and $n$ observations. The treatment is, however, quite general and could be applied to any designed experiment as outlined in the general statement given at the beginning of § $7 \cdot 1$.

Two types of question may be asked in connexion with the test for differences between groups:
(a) What is the extent of departure from the null hypothesis, measured by ( $S^{z} / \sigma_{0}^{2}$ ), that could be detected with a given chance?
(b) How many observations are we to take in each group so that we could detect a given ratio of between group to within group variability ( $S^{2} / \sigma_{0}^{2}$ ) with a prescribed chance?

To answer these questions we have to examine the function $\beta\left(S^{2} / \sigma_{0}^{2}\right)$ which may be written in the form
and consider its inverse, i.e: $\lambda \equiv \lambda\left(\nu_{1}, \nu_{2}, \alpha, \beta\right)$. Generally, $\lambda$ has to be obtained by inverse interpolation from tables of $\beta$ such as Tang's. The interval of tabulation of 0.5 for

$$
\phi=\sqrt{ }\left[\lambda /\left(\nu_{1}+1\right)\right]
$$

in Tang's tables is not fine enough for interpolation to be satisfactory. Still, they give a trial value of $\phi$ for which $\beta$ is calculated and then corrected with the help of the derivative $\partial \beta / \partial \phi$. Following this rather laborious method, Emma Lehmer (1944) has tabled $\phi$ for $\alpha=0.01$, 0.05 and $\beta=0.7,0.8$ and for a wide range of $\nu_{1}$ and $\nu_{\mathrm{g}}$. For these two values of the power we may use her tables to obtain our $\lambda$. It would clearly be of value for these tables to be extended.

We may, however, for any set of values of $\nu_{1}, \nu_{2}, \alpha$ and $\beta$, get $\lambda$ approximately with the help of the approximate form of $\beta$ given in (61). Taking a trial value of $\lambda$ we can find two consecutive integers $\lambda_{1}, \lambda_{2}$ between which $\lambda$ lies by the following method. From the expression (61) for $\beta$ we see that $\lambda$ must satisfy the relation

$$
\begin{equation*}
F_{\beta}\left(\nu, \nu_{\mathbf{a}}\right)=\frac{\nu_{1}}{\nu_{1}+\lambda} F_{a}\left(\nu_{1}, \nu_{\mathbf{a}}\right), \tag{68}
\end{equation*}
$$

where the argumente $\nu, \nu_{2}$ and $\nu_{1}, \nu_{2}$ are the degrees of freedom. Hence the two integers $\lambda_{1}$ and $\lambda_{2}$ would make the right-hand side of (68) just greater and just less than the left-hand side. These can be got by trial and error, taking the $\alpha$ and $\beta$ percentage points from the $F$-tables and comparing the two sides. (It is to be noted that $\nu$ in (66) involves $\lambda$.) For these values of $\lambda_{1}$ and $\lambda_{3}, \beta$ is then evaluated using (62) and by backward interpolation $\lambda$ is determined.

To deal with inverse problems, such as (b) mentioned above, a graphical representation of the relation between $\nu_{1}, \nu_{\mathbf{2}}$ and $\lambda$ for fixed $\alpha$ and $\beta$ will be most useful. Following the procedure described above for finding $\lambda$, charts have been constructed for $\alpha=0.05$ and for two levels of power, $\beta=0.5$ and 0.9 , which are likely to be of practical interest (see Figs. 3 (a), (b)). The charts give, to the approximation involved in (61), contours of equal power and could be used for determining any one of the three quantities, $\nu_{1}, \nu_{2}$ and $\lambda$, given the other two. When $\nu_{3}=\infty$, the $F^{\prime}$ reduces to $\chi^{\prime 2} / \nu_{1}$, and hence these charts could also be used for answering the inverse questions connected with the power function of the $\chi^{2}$-test (see p. 215).

We give here two illustrations of the use of these charts.

Illustration 1. To study the seasonal variation in the frequency of occurrence of a particular dominant alga in a pond, ten samples of 15 c.c. of water are taken from the pond on the first day of each of the five months, April to August. Fifteen drops are taken on slides from each sample after shaking it thoroughly, and the number of algae of the partioular form are


Fig. 3a. Contours of equal power for the analysis of variance test with the systematic set-up: $\alpha=0.05$, and a power $\beta\left(\nu_{1}, \nu_{2}, \lambda\right)=0.5$
counted under the microscope and the total for the fifteen slides is taken as the density for each sample.

To test whether there is significant variation in the density of this form of algae from month to month, the analysis of variance test is applied, say, at $5 \%$ level. It will be of interest to know how large should the ratio of the seasonal variability to the variability in the pond be, so that we could detect it with a $90 \%$ chance.

Here, $\nu_{1}=k-1=4, \nu_{2}=k(n-1)=4$-distributions and their applications from which we find the ratio of between . For these, the chart of Fig. $3(b)$ gives



Fig. sb. Contours of equal
This means that the odds are 9 , and a power $\beta\left(\nu_{1}, \nu_{2}, \lambda\right)=0.9$. with the syatematic set-up: of the density of the algae between to 1 on detecting differ hand, using Fig. 3 (a) wee between months was 0.58 differences at the $5 \%$ level if the s. 8.D. between months is 0.38 that there would be a 50 of the s.D. within the pond level if the s.D. Inustration 2. There are of that of a single sample in ance of detecting difference if $50: 50$ cher intended to control the are seven machines producing a month ( $S^{2} / \sigma_{0}^{\sigma}=0.145$ ). intended to control the variability in the thes producing copper pire $S^{2} / \sigma_{0}^{2}=0.145$ ).
samples from time to time and testing for differences between the machines. From previous observations we have some idea of the order of variability in the product of a single machine; suppose we do not regard the variability between machines as serious if it does not exceed 0.25 of the within-machine variability. How many samples of wire must we take from each machine to have a $90 \%$ chance of detecting, at the $5 \%$ level for $F$, a between-machine variability of this magnitude, if it exists?

Since $\frac{\lambda}{\nu_{2}}=\frac{n}{n-1} \frac{S^{2}}{\sigma_{0}^{2}}$ in virtue of (58), we have now to find $n$ satisfying the relation

$$
\frac{\lambda}{\nu_{2}}=\frac{n}{n-1} \times 0.25
$$

Following the contour in chart $3(b)$ for $\nu_{1}=6$, we find by inspection a point on it at which the ratio of the co-ordinates is nearly $0 \cdot 25$. This point gives $\nu_{2}=75$ from which we obtain the number of samples required, $n=\nu_{2} / k+1=75 / 7+1=12$, approximately. On the other hand, from $3(a)$ we find that we would have a $50 \%$ chance of detection, if $n=6$.

In conclusion, I should like to acknowledge gratefully the help and guidance I have received from Prof. E. S. Pearson and Dr H. O. Hartley in the course of my investigations.

## APPENDIX

## Distribution of the sum of squares of independent normal variates with different means and variances

Let $\xi_{1}, \xi_{2}, \ldots, \xi_{n}$ be $n$ independent normal variates with expectations $b_{1}, b_{2}, \ldots, b_{n}$ and variances $v_{1}, v_{\mathbf{g}}, \ldots, v_{n}$ respectively. The characteristic function of the statistic

$$
\begin{gather*}
\psi^{2}=\sum_{j=1}^{n} \xi_{j}^{2}  \tag{1}\\
x_{j}=\left(\xi_{j}-b_{j}\right) / \sqrt{ } v_{j}
\end{gather*}
$$

is easily obtained. Introducing
we note that each $x_{j}$ follows the probability law

$$
p(x)=\frac{1}{\sqrt{(2 \pi)}} e^{-1 x^{2}}
$$

and that $\psi^{2}$ in (1) can be written as

$$
\begin{equation*}
\psi^{2}=\Sigma v_{j}\left(x_{j}+a_{j}\right)^{2} \tag{2}
\end{equation*}
$$

where $a_{j}$ stands for $b_{j} / \sqrt{ } v_{j}$. (All summations are from $j=1$ to $n$.)
The characteristic function of $\boldsymbol{\psi}^{\mathbf{2}}$ is given by

$$
\begin{equation*}
\phi(t)=\prod_{1}^{n}\left[\frac{1}{\sqrt{(2 \pi)}} \int_{-\infty}^{\infty} \exp \left\{i t v_{j}\left(x_{j}+a_{j}\right)^{2}-\frac{1}{2} x_{j}^{8}\right\} d x_{j}\right] . \tag{3}
\end{equation*}
$$


Hence

$$
\begin{equation*}
\phi(t)=\prod_{1}^{n}\left(1-2 i t v_{j}\right)^{-i} \exp \left\{\frac{\sum\left(i t v_{j} a_{j}^{2}\right)}{1-2 i t v_{j}}\right\} \tag{4}
\end{equation*}
$$

from which all the moments of the required distribution can be derived. We may represent this approximately by a $\chi^{2}$-distribution fitted from the first two moments, $\mu_{1}^{\prime}=\Sigma v_{j}+\Sigma v_{j} a_{j}^{2}$, and $\mu_{g}=2 \Sigma v_{j}^{8}+4 \Sigma v_{j}^{8} a_{j}^{8}$.

Next we consider the conditional distribution of $\psi^{\mathbf{2}}$ in (2) subject to a single linear constraint on the $x_{j}{ }^{\prime} \mathrm{s}$, viz.

$$
\Sigma c_{j}\left(x_{j}+a_{j}\right)=\rho .
$$

The characterististic function of the joint distribution of $\psi^{\mathbf{2}}$ and $\rho$ is given by

$$
\begin{equation*}
\phi\left(t, t_{1}\right)=\Pi\left[\frac{1}{\sqrt{ }(2 \pi)} \int_{-\infty}^{\infty} \exp \left\{-\frac{1}{2} x_{j}^{2}+i t v_{j}\left(x_{j}+a_{j}\right)^{2}+i t_{1} c_{j}\left(x_{j}+a_{j}\right)\right\} d x_{j}\right] . \tag{5}
\end{equation*}
$$

On performing the integrations in (5) we find

$$
\phi\left(t, t_{1}\right)=\Pi\left[\left(1-2 i t v_{j}\right)^{-t} \exp \left\{\frac{2 i t v_{j} a_{j}^{2}+2 i t_{1} c_{j} a_{j}-t_{1}^{2} c_{j}^{2}}{2\left(1-2 i t v_{j}\right)}\right\}\right] .
$$

The conditional characteristic function of $\psi^{2}$, for fixed $\rho$ (Bartlett, 1938), is

$$
\begin{align*}
\phi(t \mid \rho)= & \int_{-\infty}^{\infty} e^{-u_{1} \rho \phi\left(t, t_{1}\right) d t_{1}} / \int_{-\infty}^{\infty} e^{-u_{1} \rho} \phi\left(0, t_{1}\right) d t_{1} \\
= & \left(\Sigma\left(\frac{c_{j}^{2}}{1-2 i t v_{j}}\right)\right)^{-i} \Pi\left(1-2 i t v_{j}\right)^{-t} \\
& \quad \times \exp \left\{\Sigma\left(\frac{i t v_{j} a_{j}^{2}}{1-2 i t v_{j}}\right)-\frac{1}{2}\left(\rho-\Sigma\left(\frac{c_{j} a_{j}}{1-2 i t v_{j}}\right)\right)^{2}\left(\Sigma\left(\frac{c_{j}^{2}}{1-2 i t v_{j}}\right)\right)^{-1}+\frac{1}{2} \frac{\left(\rho-\Sigma c_{j} a_{j}\right)^{2}}{\Sigma c_{j}^{2}}\right\} . \tag{6}
\end{align*}
$$

The moments of the conditional distribution of $\psi^{2}$ can then be obtained from (6).
Again, we may fit a Type III to the conditional distribution of $\psi^{2}$ by using the first two moments.

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[^0]:    $\dagger$ Here and below the notation $N(a, b)$ is used to indicate that a variable is normally distributed with mean $a$ and standard deviation $b$

[^1]:    $\dagger$ In making the approximation, we have associated the $\lambda$ of (37) with the distribution of $\phi^{2}$ rather than the $\lambda^{\prime}$ of (38), but this step perhaps needs fuller justification.

[^2]:    $\dagger$ [Further exploration shows that the differences between the approximate and true velues are systematic, with regular fluctustions. Use is being made of this fact to prepare certain rather more extensive tables of the power function. Ed.]

